

Relaxation approximation of some nonlinear Maxwell initial-boundary value problem

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Abstract Two nonlinear Maxwell systems are considered: Kerr model exhibiting an instantaneous response of the medium, Kerr-Debye model which contains some delay term and is a relaxation approximation of the first one. In one space dimension, we prove that the limit of the solution to the ingoing wave condition for Kerr-Debye model is a solution to the Kerr model.

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1 Introduction

Nonlinear Maxwell's equations are used for modelling nonlinear optical phenomena. The wave propagation in an isotropic medium is described by Maxwell's system:

$$\begin{aligned}\partial_t D - \operatorname{curl} H &= 0, \\ \partial_t B + \operatorname{curl} E &= 0, \\ \operatorname{div} D = \operatorname{div} B &= 0.\end{aligned}$$

The field quantities E and H represent the electric and magnetic fields, D and B the electric and magnetic displacements. We consider the constitutive relations for a nonlinear Kerr medium:

$$\begin{aligned}B &= \mu_0 H, \\ D &= \varepsilon_0 E + P,\end{aligned}$$

where P is the nonlinear polarization.

If the medium exhibits an instantaneous response we have a Kerr model:

$$P = P_K = \varepsilon_0 \varepsilon_r |E|^2 E.$$

If the medium exhibits a finite response time τ we have a Kerr-Debye model:

$$P = P_{KD} = \varepsilon_0 \chi E,$$

where

$$\partial_t \chi + \frac{1}{\tau} \chi = \frac{1}{\tau} \varepsilon_r |E|^2$$

(see for example [15] or [20]).

So the Kerr-Debye model is a relaxation approximation of the Kerr model and τ is the relaxation parameter (for a general presentation of relaxation problems, see [13]). Formally, when τ tends to 0, χ converges to $\varepsilon_r |E|^2$ and P_{KD} converges to P_K .

Recently for the Cauchy problem with regular initial data, the convergence result was obtained in [7] and [8] using the general framework given in [17]. This study is based on a precise description of the boundary layer in time.

For the modelling of realistic physical situations it is more convenient to take into account boundary value conditions. In particular we consider the impedance boundary value problem. For the Kerr-Debye model in two space dimension, numerical studies were proposed for the Dirichlet condition on the magnetic field in [19] and for the ingoing boundary condition in [8]. The aim of this paper is to prove convergence results in the case of impedance initial-boundary value problem for Kerr-Debye model in one space dimension.

2 The one space dimension models

Let us suppose that

$$\begin{aligned} D(x, y, z) &= (0, d(x), 0), \\ H(x, y, z) &= (0, 0, h(x)). \end{aligned}$$

Then Maxwell's system rewrite

$$\begin{aligned} \partial_t d + \partial_x h &= 0, \\ \partial_t h + \partial_x e &= 0, \end{aligned}$$

for $(t, x) \in ([0, +\infty[)^2$.

Once nondimensionalized the Kerr model, denoted by (K), becomes:

$$\begin{cases} \partial_t d + \partial_x h = 0, \\ \partial_t h + \partial_x e = 0, \\ d = (1 + e^2)e, \end{cases} \quad (2.1)$$

for $(t, x) \in (\mathbb{R}^+)^2$.

We suppose that the initial data vanishes

$$d(0, x) = h(0, x) = 0 \text{ for } x \in \mathbb{R}^+, \quad (2.2)$$

and that we have the boundary condition

$$h(t, 0) + ae(t, 0) = g(t) \text{ for } t \in \mathbb{R}^+, \quad (2.3)$$

where a is a non negative constant.

The system (K) is quasi-linear hyperbolic. It is a p-system where p is the reciproque function of $e \mapsto (1 + e^2)e$, and it is strictly hyperbolic with eigenvalues

$$\lambda_1 = -\sqrt{p'(d)} < 0 < \lambda_2 = \sqrt{p'(d)}.$$

The energy density \mathcal{E}_K given by

$$\mathcal{E}_K(e, h) = \frac{1}{2}(e^2 + h^2 + \frac{3}{2}e^4) \quad (2.4)$$

is a strictly convex entropy. In the entropic variable (e, h) , system (K) writes in the symmetric form:

$$\begin{pmatrix} 1 + 3e^2 & 0 \\ 0 & 1 \end{pmatrix} \partial_t \begin{pmatrix} e \\ h \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \partial_x \begin{pmatrix} e \\ h \end{pmatrix} = 0. \quad (2.5)$$

The boundary $\{x = 0\}$ is non characteristic. We can verify that the boundary condition (2.3) is an ingoing wave condition and is maximal dissipative (since $a \geq 0$). Therefore classical existence results of regular solution for the (K) boundary value problem (2.1)-(2.2)-(2.3) apply.

The one dimensional Kerr-Debye model, denoted by (KD) system, writes:

$$\begin{cases} \partial_t d_\varepsilon + \partial_x h_\varepsilon = 0, \\ \partial_t h_\varepsilon + \partial_x e_\varepsilon = 0, \\ \partial_t \chi_\varepsilon = \frac{1}{\varepsilon}(e_\varepsilon^2 - \chi_\varepsilon), \\ d_\varepsilon = (1 + \chi_\varepsilon)e_\varepsilon, \end{cases} \quad (2.6)$$

for $(t, x) \in (\mathbb{R}^+)^2$, with initial data

$$d_\varepsilon(0, x) = h_\varepsilon(0, x) = \chi_\varepsilon(0, x) = 0 \text{ for } x \in \mathbb{R}^+ \quad (2.7)$$

and with boundary condition

$$h_\varepsilon(t, 0) + ae_\varepsilon(t, 0) = g(t) \text{ for } t \in \mathbb{R}^+. \quad (2.8)$$

Using the third equation in (2.6) and (2.7) we remark that

$$\chi_\varepsilon(t, x) \geq 0. \quad (2.9)$$

The system (KD) is quasi-linear strictly hyperbolic with eigenvalues

$$\lambda_1 = (1 + \chi_\varepsilon)^{-\frac{1}{2}} < \lambda_2 = 0 < \lambda_3 = (1 + \chi_\varepsilon)^{-\frac{1}{2}}.$$

The energy density \mathcal{E}_{KD} given by

$$\mathcal{E}_{KD}(d, h, \chi) = \frac{1}{2}(1 + \chi)^{-1}d^2 + \frac{1}{2}h^2 + \frac{1}{4}\chi^2$$

is a strictly convex entropy in $\{\chi \geq 0\}$, so the system (KD) is symmetrizable.

Using the entropy variables given by

$$\begin{cases} \partial_d \mathcal{E}_{KD}(d, h, \chi) = (1 + \chi)^{-1}d := e, \\ \partial_h \mathcal{E}_{KD}(d, h, \chi) = h, \\ \partial_\chi \mathcal{E}_{KD}(d, h, \chi) = -\frac{1}{2}(1 + \chi)^{-2}d^2 + \frac{1}{2}\chi = \frac{1}{2}(\chi - e^2) := v, \end{cases}$$

system (KD) takes the symmetric form

$$\begin{cases} (1 + 3e_\varepsilon^2 + 2v_\varepsilon)\partial_t e_\varepsilon + 2e_\varepsilon\partial_t v_\varepsilon + \partial_x h_\varepsilon = 0, \\ \partial_t h_\varepsilon + \partial_x e_\varepsilon = 0, \\ 2e_\varepsilon\partial_t e_\varepsilon + 2\partial_t v_\varepsilon = -\frac{2}{\varepsilon}v_\varepsilon. \end{cases}$$

If we use the following notations:

$$W_\varepsilon = (U_\varepsilon, v_\varepsilon) \text{ with } U_\varepsilon = (e_\varepsilon, h_\varepsilon),$$

$$A_0(W_\varepsilon) = \begin{pmatrix} 1 + 3e_\varepsilon^2 + 2v_\varepsilon & 0 & 2e_\varepsilon \\ 0 & 1 & 0 \\ 2e_\varepsilon & 0 & 2 \end{pmatrix}, \quad A_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$Q(W_\varepsilon) = \begin{pmatrix} 0 \\ 0 \\ v_\varepsilon \end{pmatrix},$$

we obtain

$$A_0(W_\varepsilon)\partial_t W_\varepsilon + A_1\partial_x W_\varepsilon = -\frac{2}{\varepsilon}Q(W_\varepsilon). \quad (2.10)$$

The matrix $A_0(W_\varepsilon)$ is definite positive because $\chi_\varepsilon = e_\varepsilon^2 + 2v_\varepsilon \geq 0$ by (2.9). The matrix A_1 is singular with one dimensional kernel, so the boundary $\{x = 0\}$ is characteristic of constant multiplicity. In addition, the boundary condition (2.8) is an ingoing wave boundary condition and is maximal dissipative. Therefore the general results of [5] apply to obtain existence result of regular solutions for the (KD) boundary value problem (2.6)-(2.8).

Now let us specify the assumptions on the source term g in the boundary condition (2.3) or (2.8). For the modelling of realistic physical situations we can assume that g is compactly supported in $[0, +\infty[$. We denote by H^s the classical Sobolev space and we suppose that g belongs to $H^s(\mathbb{R})$ for s great enough. So the boundary condition and the null initial data match one each other and we obtain smooth solutions.

In order to obtain homogeneous boundary condition we replace the magnetic field h by $h(t, x) + g(t)\eta(x)$ where η is a smooth function, compactly supported in \mathbb{R}^+ , equal to 1 in a neighbourhood of 0. We rewrite the (K) boundary value problem in the entropic variables

$$\left\{ \begin{array}{l} \begin{pmatrix} 1 + 3e^2 & 0 \\ 0 & 1 \end{pmatrix} \partial_t \begin{pmatrix} e \\ h \end{pmatrix} + \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \partial_x \begin{pmatrix} e \\ h \end{pmatrix} = \begin{pmatrix} g_1 \\ g_2 \end{pmatrix} \text{ for } (t, x) \in (\mathbb{R}^+)^2, \\ e(0, x) = h(0, x) = 0 \text{ for } x \in \mathbb{R}^+, \\ h(t, 0) + ae(t, 0) = 0 \text{ for } t \in \mathbb{R}^+, \end{array} \right. \quad (2.11)$$

where $g_1(t, x) = -g(t)\eta'(x)$ and $g_2(t, x) = -g'(t)\eta(x)$, and we consider a regular solution of this problem.

Proposition 2.1 *There exists $T^* > 0$ and a unique regular solution $U = (e, h)$ for the (K) boundary value problem (2.11) defined in $[0, T^*[\times\mathbb{R}^+$ and such that*

$$\partial_t^i U \in C^0([0, T^*]; H^{2-i}(\mathbb{R}^+)) \text{ for } i = 0, 1, 2.$$

In the same way we rewrite the (KD) boundary value problem in the entropic variables

$$\left\{ \begin{array}{l} A_0(W_\varepsilon)\partial_t W_\varepsilon + A_1\partial_x W_\varepsilon = -\frac{2}{\varepsilon}Q(W_\varepsilon) + G(t, x) \text{ for } (t, x) \in (\mathbb{R}^+)^2, \\ W_\varepsilon(0, x) = 0 \text{ for } x \in \mathbb{R}^+, \\ h_\varepsilon(t, 0) + ae_\varepsilon(t, 0) = 0 \text{ for } t \in \mathbb{R}^+, \end{array} \right. \quad (2.12)$$

where $G = {}^t(g_1, g_2, 0)$.

For fixed $\varepsilon > 0$, using the results in [5] we consider a regular local solution of this problem.

Proposition 2.2 *There exists $T_\varepsilon^* > 0$ and a unique regular solution $W_\varepsilon = (U_\varepsilon, v_\varepsilon)$ for the (KD) boundary value problem (2.12) defined on $[0, T_\varepsilon^*[\times\mathbb{R}^+$ and such that*

$$\partial_t^i W_\varepsilon \in \mathcal{C}^0([0, T_\varepsilon^*]; H^{2-i}(\mathbb{R}^+)), i = 0, 1, 2.$$

The goal of this paper is to analyse the behaviour of the regular solutions for the (KD) boundary value problem when the relaxation parameter ε tends to 0. By energy estimates we first obtain uniform bound of W_ε in the following theorem.

Theorem 2.1 *Let W_ε and T_ε^* given by Proposition 2.2. There exist $\tilde{T} > 0$ and a constant $K > 0$ such that for all $\varepsilon > 0$, $T_\varepsilon^* \geq \tilde{T}$ and the solution $W_\varepsilon = (U_\varepsilon, v_\varepsilon)$ of the (KD) boundary value problem (2.12) satisfies*

$$\left\{ \begin{array}{l} \|\partial_t^i U_\varepsilon\|_{\mathcal{C}^0([0, \tilde{T}]; H^{2-i}(\mathbb{R}^+))} \leq K \text{ for } i = 0, 1, 2, \\ \frac{1}{\varepsilon} \|v_\varepsilon\|_{\mathcal{C}^0([0, \tilde{T}]; H^1(\mathbb{R}^+))} \leq K, \\ \|\partial_t^i v_\varepsilon\|_{\mathcal{C}^0([0, \tilde{T}]; H^{2-i}(\mathbb{R}^+))} \leq K \text{ for } i = 0, 1, 2. \end{array} \right.$$

Theorem 2.1 shows the strong convergence of $v_\varepsilon = \frac{1}{2}(\chi_\varepsilon - e_\varepsilon^2)$ to zero. The convergence of U_ε to the solution of the boundary value problem (2.11) for (K) is contained in the following statement.

Theorem 2.2 *Let U and T^* given by Proposition 2.1, let U_ε and \tilde{T} given by Theorem 2.1. For $T \leq \tilde{T}$ and $T < T^*$, there exists a constant $K > 0$ such that for all $\varepsilon > 0$,*

$$\|U_\varepsilon - U\|_{\mathcal{C}^0([0, T]; L^2(\mathbb{R}^+))} \leq K\varepsilon. \quad (2.13)$$

So we have proved that the regular solutions for the (KD) boundary value problem tend to the solution for the (K) boundary value problem when the relaxation parameter tends to zero.

In our study we remark that no boundary layer appears in the time variable because the null initial data belongs to the equilibrium manifold defined by

$$\mathcal{V} = \{(d, h, \chi) \text{ such that } \chi - (1 + \chi)^{-2}d^2 = 2v = 0\}.$$

For the space variable, we have the same boundary condition for the system (K) and for the system (KD), so no space boundary layer appears again.

In the case of non characteristic boundary conditions, a general study of boundary conditions for hyperbolic relaxation systems is given in [16].

To our knowledge general convergence results are not available for hyperbolic relaxation systems in domains with boundary in the literature.

A special well investigated problem is the semilinear relaxation approximation to the boundary value problem for a scalar quasilinear equation, see [11, 14, 10, 12], and [4, 1] for related numerical considerations.

For the strong solutions of the (KD) boundary value problem with the entropic variables, we obtain a symmetric hyperbolic system endowed with a flat equilibrium manifold. This basic structure of numerous relaxation systems is explained and used for the global existence of smooth solutions in [6] and for asymptotic behaviour in [2]. For connected works see also [18, 9].

The study of the three-dimensional case is the subject of a work in progress [3]. In this case the previous properties are still valid: there is no boundary layer and with the entropic variables, the equilibrium manifold is flat. On the other hand, the boundary is characteristic for both problems, with a two-dimensional kernel for Kerr problem, and a three-dimensional kernel for Kerr-Debye problem. We must then take into account the nonlinear conservation equations $\operatorname{div} D = \operatorname{div} B = 0$, which is irrelevant for the one-dimensional case.

Section 3 is devoted to the proof of Theorem 2.1. Using energy estimates, we bound the time derivatives of W_ε . The boundary is characteristic so we cannot directly obtain the bounds for the space derivatives. Since $\operatorname{Ker} A_1 = \mathbb{R}(0, 0, 1)$ with the first two equations, we can express and estimate the space derivatives for U_ε . Estimates for v_ε are obtained solving the third equation by Duhamel formula.

We prove Theorem 2.2 in Section 4.

3 Proof of Theorem 2.1

We denote by \cdot the canonical scalar product on \mathbb{R}^3 and by $|\cdot|$ the associated euclidean norm.

For the convenience of the reader, we rewrite System (2.12) omitting the dependance on ε . Let W be a solution to the problem

$$\left\{ \begin{array}{l} A_0(W)\partial_t W + A_1\partial_x W = -\frac{2}{\varepsilon}Q(W) + G \text{ on } [0, T_\varepsilon^*[\times\mathbb{R}^+, \\ W(0, x) = 0 \text{ on } \mathbb{R}^+, \\ (h + ae)(t, 0) = 0 \text{ on } [0, T_\varepsilon^*[. \end{array} \right. \quad (3.1)$$

Here we denote by $T_\varepsilon^* \in]0, +\infty]$ the lifespan of W given by [5]. We prove the result when

$$\|W\|_{L^\infty(0, T_\varepsilon^* \times \mathbb{R}^+)} = +\infty.$$

If this assumption is not satisfied, we use the extension Theorem in [5] with analogous arguments as below.

Let M be a positive constant large enough with respect to the L^∞ -norm of the initial data. Consider the regular solution $W = (U, v)$ to problem (3.1), as given by Proposition 2.2. Define $T_\varepsilon > 0$ as

$$T_\varepsilon = \max\{T \leq T_\varepsilon^*, \|W\|_{L^\infty([0, T] \times \mathbb{R}^+)} < M\}, \quad (3.2)$$

that is T_ε is the first time such that the L^∞ -norm of W reaches the given bound M .

We recall that the solution $W = (U, v)$ satisfies:

$$\partial_t^i W \in \mathcal{C}^0([0, T_\varepsilon]; H^{2-i}(\mathbb{R}^+)) \text{ for } i = 0, 1, 2. \quad (3.3)$$

Furthermore we recall that

$$e^2 + 2v \geq 0 \text{ on } [0, T_\varepsilon] \times \mathbb{R}^+. \quad (3.4)$$

From (3.2) and (3.4), by definition of A_0 we have

$$\forall \xi \in \mathbb{R}^3, A_0(W)\xi \cdot \xi \geq \frac{1}{M^2 + 2} |\xi|^2. \quad (3.5)$$

We measure the boundary value lifting G by the quantities γ and Γ defined by:

$$\gamma(t) = \left(\|G(t)\|_{L^2(\mathbb{R}^+)}^2 + \|\partial_t G(t)\|_{L^2(\mathbb{R}^+)}^2 + \|\partial_{tt} G(t)\|_{L^2(\mathbb{R}^+)}^2 + \|\partial_x G(t)\|_{L^2(\mathbb{R}^+)}^2 \right)^{\frac{1}{2}}, \quad (3.6)$$

and

$$\Gamma(t) = \sup_{s \in [0, t]} \gamma(s). \quad (3.7)$$

For $t \in [0, T_\varepsilon]$ we define φ and Φ by

$$\varphi(t) = \left(\|W(t)\|_{L^2(\mathbb{R}^+)}^2 + \|\partial_t W(t)\|_{L^2(\mathbb{R}^+)}^2 + \|\partial_{tt} W(t)\|_{L^2(\mathbb{R}^+)}^2 \right)^{\frac{1}{2}}, \quad (3.8)$$

$$\Phi(t) = \sup_{s \in [0, t]} \varphi(s). \quad (3.9)$$

We first prove the following result.

Lemma 3.1 *There exists a constant K_1 , independant of M , such that*

$$\forall \varepsilon > 0, \forall t \in [0, T_\varepsilon], \Phi(t)^2 \leq K_1(2 + M^2) \int_0^t \left(1 + (\Phi(s))^5 + (\Gamma(s))^5 \right) ds. \quad (3.10)$$

First step: L^2 estimate

We take the inner product of Equation (3.1) with W and we obtain that

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W)W \cdot W dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |v|^2 dx + a|e(t, 0)|^2 = \int_{\mathbb{R}^+} G \cdot W + \frac{1}{2} \int_{\mathbb{R}^+} \partial_t(A_0(W))W \cdot W dx,$$

where

$$\partial_t(A_0(W)) = \begin{pmatrix} 6e\partial_t e + 2\partial_t v & 0 & 2\partial_t e \\ 0 & 0 & 0 \\ 2\partial_t e & 0 & 0 \end{pmatrix}. \quad (3.11)$$

We have

$$\begin{aligned} \left| \int_{\mathbb{R}^+} \partial_t(A_0(W))W \cdot W dx \right| &\leq \|\partial_t(A_0(W))\|_{L^\infty(\mathbb{R}^+)} \|W\|_{L^2(\mathbb{R}^+)}^2 \\ &\leq C(1 + \|e\|_{L^\infty(\mathbb{R}^+)}) (\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v\|_{L^\infty(\mathbb{R}^+)}) \|W\|_{L^2(\mathbb{R}^+)}^2, \end{aligned}$$

thus we obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W) W \cdot W \, dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |v|^2 \, dx + a|e(t, 0)|^2 &\leq \|G\|_{L^2(\mathbb{R}^+)} \|W\|_{L^2(\mathbb{R}^+)} \\ &+ C(1 + \|e\|_{L^\infty(\mathbb{R}^+)}) (\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v\|_{L^\infty(\mathbb{R}^+)}) \|W\|_{L^2(\mathbb{R}^+)}^2. \end{aligned} \quad (3.12)$$

Second step: estimate on $\partial_t W$

We can derivate the system (3.1) with respect to t . We obtain that

$$\begin{cases} A_0(W) \partial_{tt} W + \partial_t(A_0(W)) \partial_t W + A_1 \partial_{tx} W = -\frac{2}{\varepsilon} Q(\partial_t W) + \partial_t G \text{ on } [0, T_\varepsilon] \times \mathbb{R}^+, \\ \partial_t W(0, x) = 0 \text{ on } \mathbb{R}^+, \\ (\partial_t h + a \partial_t e)(t, 0) = 0 \text{ on } [0, T_\varepsilon]. \end{cases} \quad (3.13)$$

Taking the inner product of (3.13) with $\partial_t W$, we obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W) \partial_t W \cdot \partial_t W \, dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |\partial_t v|^2 \, dx + a|\partial_t e(t, 0)|^2 &= \int_{\mathbb{R}^+} \partial_t G \cdot \partial_t W \\ &- \frac{1}{2} \int_{\mathbb{R}^+} \partial_t(A_0(W)) \partial_t W \cdot \partial_t W \, dx, \end{aligned}$$

and thus there exists a constant C such that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W) \partial_t W \cdot \partial_t W \, dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |\partial_t v|^2 \, dx + a|\partial_t e(t, 0)|^2 &\leq \|\partial_t G\|_{L^2(\mathbb{R}^+)} \|\partial_t W\|_{L^2(\mathbb{R}^+)} \\ &+ C(1 + \|e\|_{L^\infty(\mathbb{R}^+)}) (\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v\|_{L^\infty(\mathbb{R}^+)}) \|\partial_t W\|_{L^2(\mathbb{R}^+)}^2. \end{aligned} \quad (3.14)$$

Third step: estimate on $\partial_{tt} W$

We can derivate System (3.13) with respect to t . We obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W) \partial_{tt} W \cdot \partial_{tt} W \, dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |\partial_{tt} v|^2 \, dx + a|\partial_{tt} e(t, 0)|^2 &= \int_{\mathbb{R}^+} \partial_{tt} G \cdot \partial_{tt} W \\ &- \frac{3}{2} \int_{\mathbb{R}^+} \partial_t(A_0(W)) \partial_{tt} W \cdot \partial_{tt} W \, dx - \int_{\mathbb{R}^+} \partial_{tt}(A_0(W)) \partial_t W \cdot \partial_{tt} W. \end{aligned}$$

Now we have:

$$\partial_{tt}(A_0(W)) = \begin{pmatrix} 6(\partial_t e)^2 + 6e \partial_{tt} e + 2\partial_{tt} v & 0 & 2\partial_{tt} e \\ 0 & 0 & 0 \\ 2\partial_{tt} e & 0 & 0 \end{pmatrix},$$

thus

$$\begin{aligned} \|\partial_{tt}(A_0(W)) \partial_t W\|_{L^2(\mathbb{R}^+)} &\leq C(1 + \|e\|_{L^\infty(\mathbb{R}^+)}) (\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v\|_{L^\infty(\mathbb{R}^+)}) \|\partial_{tt} W\|_{L^2(\mathbb{R}^+)} \\ &+ C \|\partial_t e\|_{L^\infty(\mathbb{R}^+)}^2 \|\partial_t e\|_{L^2(\mathbb{R}^+)}. \end{aligned}$$

Therefore

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} A_0(W) \partial_{tt} W \cdot \partial_{tt} W dx + \frac{2}{\varepsilon} \int_{\mathbb{R}^+} |\partial_{tt} v|^2 dx + a |\partial_{tt} e(t, 0)|^2 \\
& \leq \|\partial_{tt} G\|_{L^2(\mathbb{R}^+)} \|\partial_{tt} W\|_{L^2(\mathbb{R}^+)} + C(1 + \|e\|_{L^\infty(\mathbb{R}^+)}) (\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v\|_{L^\infty(\mathbb{R}^+)}) \|\partial_{tt} W\|_{L^2(\mathbb{R}^+)}^2 \\
& \quad + C \|\partial_t e\|_{L^\infty(\mathbb{R}^+)}^2 \|\partial_t e\|_{L^2(\mathbb{R}^+)} \|\partial_{tt} W\|_{L^2(\mathbb{R}^+)}.
\end{aligned} \tag{3.15}$$

Fourth step: L^∞ estimates for $\partial_t e$ and $\partial_t v$

We recall the equations satisfied by $W = (e, h, v)$:

$$(3e^2 + 2v + 1)\partial_t e + 2e\partial_t v + \partial_x h = g_1, \tag{3.16}$$

$$\partial_t h + \partial_x e = g_2, \tag{3.17}$$

$$e\partial_t e + \partial_t v = -\frac{1}{\varepsilon}v. \tag{3.18}$$

From (3.17) we have

$$\|\partial_x e\|_{L^2(\mathbb{R}^+)} \leq \varphi + \|G\|_{L^2(\mathbb{R}^+)} \leq \varphi + \gamma, \tag{3.19}$$

where φ is defined by (3.8) and γ is defined by (3.6). Thus

$$\|e\|_{H^1(\mathbb{R}^+)} \leq 2\varphi + \gamma, \tag{3.20}$$

and by Sobolev injection

$$\|e\|_{L^\infty(\mathbb{R}^+)} \leq C(\varphi + \gamma). \tag{3.21}$$

Derivating (3.17) with respect to t we obtain that

$$\|\partial_x \partial_t e\|_{L^2(\mathbb{R}^+)} \leq \varphi + \|\partial_t G\|_{L^2(\mathbb{R}^+)},$$

so

$$\|\partial_t e\|_{H^1(\mathbb{R}^+)} \leq 2\varphi + \|\partial_t G\|_{L^2(\mathbb{R}^+)}, \tag{3.22}$$

and by Sobolev injections there exists a constant C such that

$$\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} \leq C(\varphi + \gamma). \tag{3.23}$$

Now we solve Equation (3.18) with Duhamel formula:

$$v(t, x) = - \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) e(s, x) \partial_t e(s, x) ds,$$

so

$$\begin{aligned}
\|v(t, \cdot)\|_{L^\infty(\mathbb{R}^+)} & \leq \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) \|e(s, \cdot)\|_{L^\infty(\mathbb{R}^+)} \|\partial_t e(s, \cdot)\|_{L^\infty(\mathbb{R}^+)} ds \\
& \leq C \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) (\varphi(s) + \gamma(s))^2 ds \\
& \leq C(\Phi(t) + \Gamma(t))^2 \varepsilon,
\end{aligned}$$

where Φ is defined in (3.9) and Γ is defined by (3.7).

Using Equation (3.18) we obtain then that

$$\begin{aligned} \|\partial_t v\|_{L^\infty(\mathbb{R}^+)} &\leq \frac{1}{\varepsilon} \|v\|_{L^\infty(\mathbb{R}^+)} + \|e\partial_t e\|_{L^\infty(\mathbb{R}^+)} \\ &\leq C(\Phi(t) + \Gamma(t))^2. \end{aligned} \quad (3.24)$$

Fifth step: end of the proof of Lemma 3.1

We sum up inequalities (3.12), (3.14) and (3.15). Using (3.23) and (3.24) we obtain that there exists a constant K_1 independent of ε , M and $t \in [0, T_\varepsilon]$ such that

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} \left(A_0(W)W \cdot W dx + A_0(W)\partial_t W \cdot \partial_t W dx + A_0(W)\partial_{tt} W \cdot \partial_{tt} W \right) dx \\ &+ \frac{2}{\varepsilon} \int_{\mathbb{R}^+} \left(|v|^2 dx + |\partial_t v|^2 dx + |\partial_{tt} v|^2 \right) dx + a \left(|e(t, 0)|^2 + |\partial_t e(t, 0)|^2 + |\partial_{tt} e(t, 0)|^2 \right) \\ &\leq C\gamma\varphi + C(1 + \varphi + \gamma) (\varphi + \gamma + (\Phi + \Gamma)^2) \varphi^2 + C(\varphi + \gamma)^2 \varphi^2 \\ &\leq K_1 (1 + \Phi^5 + \Gamma^5). \end{aligned}$$

We integrate this inequality on $[0, t]$ for $t \in [0, T_\varepsilon]$ and we have by (3.5) that

$$\forall t \in [0, T_\varepsilon], \varphi^2(t) \leq K_1(2 + M^2) \int_0^t (1 + \Phi(s)^5 + \Gamma(s)^5) ds.$$

So we obtain Lemma 3.1.

Lemma 3.1 provides estimates on the time derivatives of W . The space derivatives can not be obtained by the same method since it is impossible to derivate the system (KD) with respect to the normal variable x . We deduce estimates on the space derivatives from the equations as we will see in the following lemma.

Lemma 3.2 *There exists K_2 such that for all $t < T_\varepsilon^*$,*

$$\|U\|_{H^2(\mathbb{R}^+)} \leq K_2(\Gamma + \Phi + \Gamma^3 + \Phi^3),$$

$$\|v\|_{H^1(\mathbb{R}^+)} \leq K_2(\Gamma^2 + \Phi^2)\varepsilon,$$

$$\|\partial_t v\|_{H^1(\mathbb{R}^+)} \leq K_2(\Gamma^2 + \Phi^2),$$

$$\|W\|_{L^\infty(\mathbb{R}^+)} \leq K_2(\Gamma + \Phi + \Gamma^3 + \Phi^3).$$

Proof of Lemma 3.2

We recall that

$$v(t, x) = - \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) e(s, x) \partial_t e(s, x) ds,$$

thus

$$\|v(t, \cdot)\|_{H^1(\mathbb{R}^+)} \leq C \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) \|e(s, \cdot)\|_{H^1(\mathbb{R}^+)} \|\partial_t e(s, \cdot)\|_{H^1(\mathbb{R}^+)} ds,$$

thus by Estimates (3.22) and (3.20), we have

$$\|v(t, \cdot)\|_{H^1(\mathbb{R}^+)} \leq C(\Phi(t) + \Gamma(t))^2 \varepsilon. \quad (3.25)$$

From (3.18) we obtain that

$$\|\partial_t v(t, \cdot)\|_{H^1(\mathbb{R}^+)} \leq \frac{1}{\varepsilon} \|v(t, \cdot)\|_{H^1(\mathbb{R}^+)} + \|e\|_{H^1(\mathbb{R}^+)} \|\partial_t e\|_{H^1(\mathbb{R}^+)},$$

thus

$$\|\partial_t v(t, \cdot)\|_{H^1(\mathbb{R}^+)} \leq C(\Phi(t) + \Gamma(t))^2. \quad (3.26)$$

Now from Equation (3.16) we estimate $\partial_x h$ in the following way:

$$\|\partial_x h\|_{H^1(\mathbb{R}^+)} \leq \|g_1\|_{H^1(\mathbb{R}^+)} + \|3e^2 + 2v\|_{H^1(\mathbb{R}^+)} \|\partial_t e\|_{H^1(\mathbb{R}^+)} + \|\partial_t e\|_{H^1(\mathbb{R}^+)} + 2\|e\|_{H^1(\mathbb{R}^+)} \|\partial_t v\|_{H^1(\mathbb{R}^+)},$$

and thus there exists a constant C such that

$$\|h\|_{H^2(\mathbb{R}^+)} \leq C(\Gamma + \Phi + \Gamma^3 + \Phi^3). \quad (3.27)$$

We derivate (3.17) with respect to x :

$$\partial_{xx} e = \partial_x g_2 - \partial_t \partial_x h.$$

Thus

$$\|\partial_{xx} e\|_{L^2(\mathbb{R}^+)} = \gamma + \|\partial_t \partial_x h\|_{L^2(\mathbb{R}^+)}.$$

Derivating (3.16) with respect to t we obtain that:

$$\begin{aligned} \|\partial_t \partial_x h\|_{L^2(\mathbb{R}^+)} &\leq \|3e^2 + 2v + 1\|_{L^\infty(\mathbb{R}^+)} \|\partial_{tt} e\|_{L^2(\mathbb{R}^+)} + \|6e\partial_t e + 2\partial_t v\|_{L^2(\mathbb{R}^+)} \|\partial_t e\|_{L^\infty(\mathbb{R}^+)} \\ &\quad + 2\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} \|\partial_t v\|_{L^2(\mathbb{R}^+)} + 2\|\partial_t e\|_{L^\infty(\mathbb{R}^+)} \|\partial_{tt} v\|_{L^2(\mathbb{R}^+)} + \|\partial_t g_1\|_{L^2(\mathbb{R}^+)} \\ &\leq C(\Gamma + \Phi + \Phi^3 + \Gamma^3), \end{aligned}$$

using (3.21), (3.23), (3.24) and (3.25).

So we have obtained that

$$\|e\|_{H^2(\mathbb{R}^+)} \leq C(\Gamma + \Phi + \Phi^3 + \Gamma^3). \quad (3.28)$$

This concludes the proof of Lemma 3.2.

Proof of Theorem 2.1

We fix $M = 3K_2(\Gamma_\infty + \Gamma_\infty^3)$, where K_2 is given by Lemma 3.2, and where $\Gamma_\infty = \sup_{\mathbb{R}^+} \Gamma$.

Let us introduce ξ the solution of

$$\begin{cases} \xi' = K_1(2 + M^2) \left(1 + \xi^{\frac{5}{2}} + \Gamma_\infty^5\right), \\ \xi(0) = 0, \end{cases}$$

defined on the maximal interval $[0, T_1[$, where K_1 is given by Lemma 3.1. By comparison results, from Estimate (3.10) we have

$$\forall t \in [0, \min(T_1, T_\varepsilon)[, \quad \Phi^2(t) \leq \xi(t).$$

From Lemma 3.2 we have then that

$$\|W\|_{L^\infty(\mathbb{R}^+)} \leq K_2(\Gamma + \xi^{\frac{1}{2}} + \Gamma^3 + \xi^{\frac{3}{2}}).$$

The map ξ is continuous and $\xi(0) = 0$, thus there exists a time \tilde{T} with $0 < \tilde{T} \leq T_1$ such that

$$K_2(\xi^{\frac{1}{2}}(\tilde{T}) + \xi^{\frac{3}{2}}(\tilde{T})) \leq \frac{M}{3}.$$

For all $t \leq \tilde{T}$, we then have:

$$K_2(\Gamma + \xi^{\frac{1}{2}} + \Gamma^3 + \xi^{\frac{3}{2}}) \leq \frac{2M}{3}.$$

Thus $\tilde{T} < T_\varepsilon$ because if $T_\varepsilon \leq \tilde{T}$ then $\|W(t)\|_{L^\infty(\mathbb{R}^+)} \leq \frac{2M}{3}$ for all $t \leq T_\varepsilon$, which is contradictory to the fact that $\|W(T_\varepsilon)\|_{L^\infty(\mathbb{R}^+)} = M$ (see (3.2)).

It remains to obtain an uniform H^2 estimate on v . We have

$$v(t, x) = -\frac{1}{2} \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) \partial_t(e(s, x)^2) ds,$$

thus with an integration by parts we obtain that

$$v(t, x) = -\frac{1}{2} \left[\exp\left(\frac{s-t}{\varepsilon}\right) e(s, x)^2 \right]_0^t + \frac{1}{2\varepsilon} \int_0^t \exp\left(\frac{s-t}{\varepsilon}\right) e(s, x)^2 ds.$$

So we obtain that there exists a constant C independant of ε such that for $t \leq \tilde{T}$

$$\|v(t, \cdot)\|_{H^2(\mathbb{R}^+)} \leq C \|e\|_{L^\infty([0, t]; H^2(\mathbb{R}^+))}^2.$$

This concludes the proof of Theorem 2.1.

4 Proof of Theorem 2.2

Let us consider $U = (e, h)$ the solution of the (K) boundary value problem (2.11) defined on the time interval $[0, T^*[$ given by Proposition 2.1. We recall that we denote by $W_\varepsilon = (U_\varepsilon, v_\varepsilon)$ the solution of the (KD) boundary value problem (2.12) given by Proposition 2.2, with $U_\varepsilon = (e_\varepsilon, h_\varepsilon)$.

Set $R_\varepsilon = \frac{1}{\varepsilon}(U_\varepsilon - U) = (r_\varepsilon, s_\varepsilon)$. The remainder term R_ε satisfies

$$\begin{cases} (3e_\varepsilon^2 + 2v_\varepsilon + 1)\partial_t r_\varepsilon + \partial_x s_\varepsilon + 3(e_\varepsilon + e)(\partial_t e)r_\varepsilon = -\frac{2}{\varepsilon}e_\varepsilon\partial_t v_\varepsilon - \frac{2}{\varepsilon}v_\varepsilon\partial_t e, \\ \partial_t s_\varepsilon + \partial_x r_\varepsilon = 0, \end{cases} \quad (4.1)$$

with the initial and boundary conditions:

$$\begin{cases} s_\varepsilon = e_\varepsilon = 0 \text{ for } t = 0, \\ s_\varepsilon + ar_\varepsilon = 0 \text{ for } x = 0. \end{cases} \quad (4.2)$$

In order to estimate the right hand side term we recall that from (2.12) we have

$$\frac{1}{\varepsilon}\partial_t v_\varepsilon = -\partial_{tt}v_\varepsilon - (\partial_t e_\varepsilon)^2 - e_\varepsilon\partial_{tt}e_\varepsilon.$$

Using Theorem 2.1, for all $T \leq \tilde{T}$, there exists a constant K independent of ε such that

$$\|\partial_t v_\varepsilon\|_{L^2(\mathbb{R}^+)} \leq K\varepsilon. \quad (4.3)$$

We take the inner product of (4.1) with R_ε to obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} ((3e_\varepsilon^2 + 2v_\varepsilon + 1)r_\varepsilon^2 + s_\varepsilon^2) + a(s_\varepsilon(t, 0))^2 &= \frac{1}{2} \int_{\mathbb{R}^+} (6e_\varepsilon \partial_t e_\varepsilon + 2\partial_t v_\varepsilon) r_\varepsilon^2 \\ &- \int_{\mathbb{R}^+} 3(e_\varepsilon + e)(\partial_t e) r_\varepsilon^2 - \frac{2}{\varepsilon} \int_{\mathbb{R}^+} e_\varepsilon (\partial_t v_\varepsilon) r_\varepsilon - \frac{2}{\varepsilon} \int_{\mathbb{R}^+} v_\varepsilon (\partial_t e) r_\varepsilon. \end{aligned}$$

So, there exists a constant C such that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} ((3e_\varepsilon^2 + 2v_\varepsilon + 1)r_\varepsilon^2 + s_\varepsilon^2) &\leq C (\|e_\varepsilon\|_{L^\infty(\mathbb{R}^+)} \|\partial_t e_\varepsilon\|_{L^\infty(\mathbb{R}^+)} + \|\partial_t v_\varepsilon\|_{L^\infty(\mathbb{R}^+)}) \|r_\varepsilon\|_{L^2(\mathbb{R}^+)}^2 \\ &+ C (\|e_\varepsilon\|_{L^\infty(\mathbb{R}^+)} \|\partial_t e\|_{L^\infty(\mathbb{R}^+)} + \|e\|_{L^\infty(\mathbb{R}^+)} \|\partial_t e\|_{L^\infty(\mathbb{R}^+)}) \|r_\varepsilon\|_{L^2(\mathbb{R}^+)}^2 \\ &+ \frac{2}{\varepsilon} \|e_\varepsilon\|_{L^\infty(\mathbb{R}^+)} \|\partial_t v_\varepsilon\|_{L^2(\mathbb{R}^+)} \|r_\varepsilon\|_{L^2(\mathbb{R}^+)} + \frac{2}{\varepsilon} \|\partial_t e\|_{L^\infty(\mathbb{R}^+)} \|v_\varepsilon\|_{L^2(\mathbb{R}^+)} \|r_\varepsilon\|_{L^2(\mathbb{R}^+)}. \end{aligned}$$

From the estimates of Theorem 2.1, from the properties of e (see Proposition 2.1) and from (4.3), there exists a constant K such that for all $T \leq \min(\tilde{T}, T^*)$,

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^+} ((3e_\varepsilon^2 + 2v_\varepsilon + 1)r_\varepsilon^2 + s_\varepsilon^2) \leq K(1 + \|r_\varepsilon\|_{L^2(\mathbb{R}^+)}^2).$$

We integrate this equation from 0 to t and we obtain, using (3.5), that

$$\|r_\varepsilon(t)\|_{L^2(\mathbb{R}^+)}^2 + \|s_\varepsilon(t)\|_{L^2(\mathbb{R}^+)}^2 \leq K(M^2 + 2) \int_0^t (1 + \|r_\varepsilon(\tau)\|_{L^2(\mathbb{R}^+)}^2) d\tau.$$

We conclude the proof of Theorem 2.2 by Gronwall Lemma.

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