

# Optimized transmission conditions for domain decomposition methods and Helmholtz equation. Application to higher order finite element methods.

F. Collino, M. Duruflé, M. Lecouvez, P. Joly

23 June 2014

# Model problem

Helmholtz equation

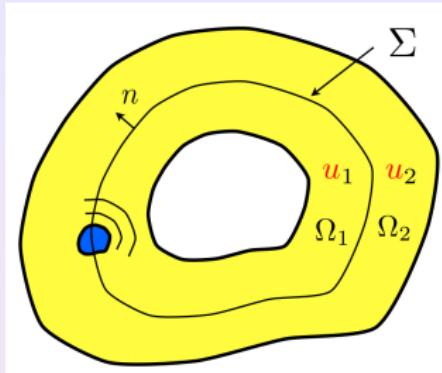
$$\begin{cases} -\rho \omega^2 \mathbf{u} - \operatorname{div}(\mu \nabla \mathbf{u}) = 0, & \text{in } \Omega \\ \mathbf{u} = 0, & \text{on } \Gamma_1 \\ \frac{\partial \mathbf{u}}{\partial n} - i k(\omega) \mathbf{u} = \frac{\partial \mathbf{u}^{\text{inc}}}{\partial n} - i k(\omega) \mathbf{u}^{\text{inc}}, & \text{on } \Gamma_2 \end{cases}$$

with  $k(\omega)$  the wave number :

$$k(\omega) = \omega \sqrt{\frac{\rho}{\mu}}$$

and  $\mathbf{u}^{\text{inc}}$  an incident plane wave.

# Transmission conditions



Transmission conditions :

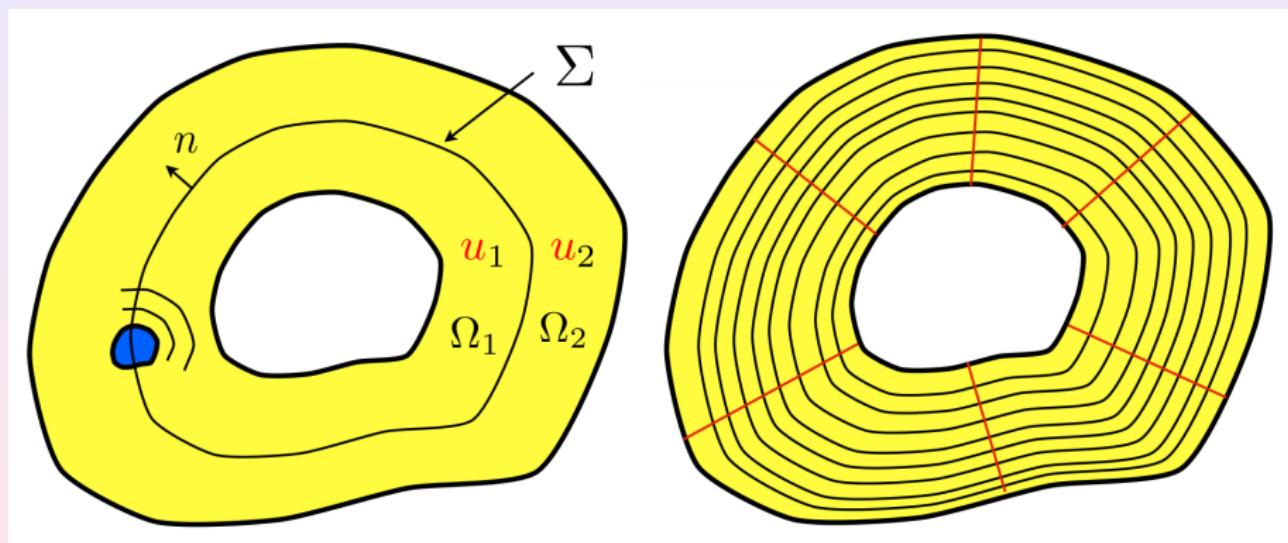
$$u_1 = u_2, \quad \mu_1 \partial_n u_1 = \mu_2 \partial_n u_2$$

Iterative DDMs : produce a sequence  $(u_1^n, u_2^n)$

- $(u_1^n, u_2^n)$  computed from previous iterations by solving local problems in  $\Omega_1$  and  $\Omega_2$
- $(u_1^n, u_2^n) \rightarrow (u_1, u_2)$  when  $n$  tends to the infinity

# Transmission conditions

Concentric interfaces without intersection:



# Equivalent transmission conditions

Transmission conditions are rewritten (coefficients  $\rho, \mu$  continuous across the interface  $\Sigma$ )

$$\mu(\partial_n u_1 - z [k(\omega)u_1 + cT(u_1)]) = \mu(\partial_n u_2 - z [k(\omega)u_2 + cT(u_2)])$$

$$\mu(\partial_n u_1 - \bar{z} [k(\omega)u_1 + cT(u_1))] = \mu(\partial_n u_2 - \bar{z} [k(\omega)u_2 + cT(u_2)])$$

Let us denote

$$B_{z,c} = \mu(\partial_n - z [k(\omega)\mathbb{I} + cT])$$

Transmission conditions given as :

$$B_{z,c}u_1 = B_{z,c}u_2, \quad B_{\bar{z},c}u_1 = B_{\bar{z},c}u_2$$

# Jacobi iterative algorithm

Sequence  $(\mathbf{u}_1^n, \mathbf{u}_2^n)$  obtained with Jacobi iterative algorithm:

$$\begin{cases} -\rho \omega^2 \mathbf{u}_1^n - \operatorname{div}(\mu \mathbf{u}_1^n) = 0 & \text{in } \Omega_1 \\ -\rho \omega^2 \mathbf{u}_2^n - \operatorname{div}(\mu \mathbf{u}_2^n) = 0 & \text{in } \Omega_2 \\ \mathcal{B}_{z,c} \mathbf{u}_1^n = \mathcal{B}_{z,c} \mathbf{u}_2^{n-1} \\ \mathcal{B}_{\bar{z},c} \mathbf{u}_2^n = \mathcal{B}_{\bar{z},c} \mathbf{u}_1^{n-1} \end{cases}$$

Relaxation with parameter  $r$ :

$$\mathcal{B}_{z,c} \mathbf{u}_1^n = r \mathcal{B}_{z,c} \mathbf{u}_2^{n-1} + (1 - r) \mathcal{B}_{z,c} \mathbf{u}_1^{n-1}$$

$$\mathcal{B}_{\bar{z},c} \mathbf{u}_2^n = r \mathcal{B}_{\bar{z},c} \mathbf{u}_1^{n-1} + (1 - r) \mathcal{B}_{\bar{z},c} \mathbf{u}_2^{n-1}$$

# Classical choices for operator B

Scalar Impedance

$$z = i \quad T = 0$$

- Després (1990), Després, Joly, Roberts (1992)
- Collino, Ghanemi, Joly (1998)

Local operators :

$$z = i \quad T = (\mathbb{I} - \alpha_1 \Delta_\Sigma)^{-1} (\mathbb{I} - \alpha_2 \Delta_\Sigma)$$

- Gander, Magoules, Nataf (2002), Japhet, Nataf (2002)
- J.F. Lee (2006), Antoine, Boubendir, Geuzaine (2012)

# Exponential convergence

For a non-local operator  $\textcolor{blue}{T}$  of the form

$$\textcolor{blue}{T} = \Lambda \Lambda^*$$

where  $\Lambda$  is an isomorphism from  $L^2(\Sigma)$  to  $H^{-1/2}(\Sigma)$ , there exists  $\tau(\textcolor{magenta}{r}, \textcolor{violet}{z}, \textcolor{red}{c}, \Lambda)$  such that

$$\|\textcolor{red}{u}_1^n - \textcolor{red}{u}_1\| + \|\textcolor{red}{u}_2^n - \textcolor{red}{u}_2\| \leq C \tau^n$$

Optimization of parameters  $\textcolor{violet}{z}$ ,  $\textcolor{red}{c}$  for circular layers.

# Class of non-local operators

$\Lambda$  pseudo-differential operator of order 1/2, its symbol would be

$$\hat{\Lambda} = \left| \frac{\xi}{\omega} \right|^{1/2}$$

In 2-D, an operator satisfying these properties is :

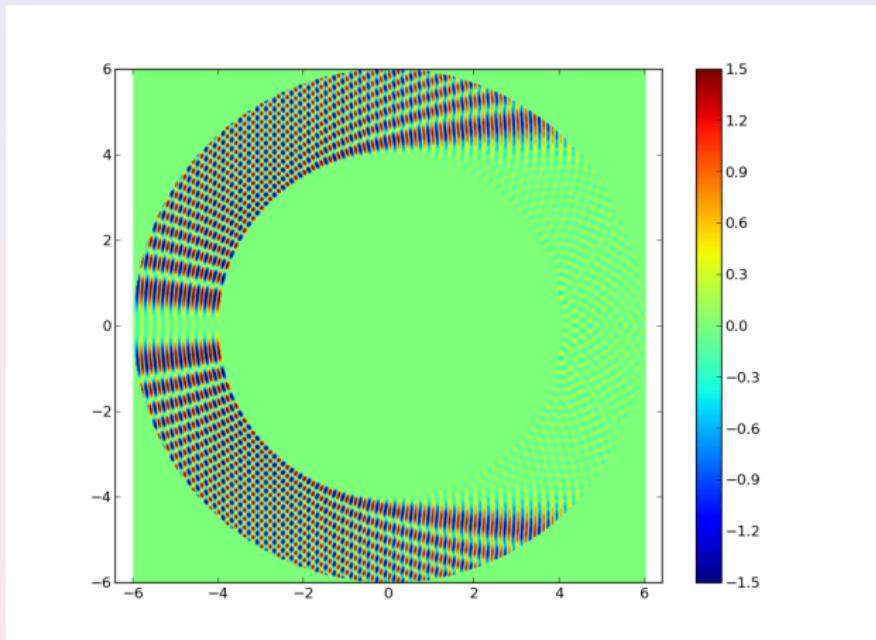
$$(\Lambda u, \varphi) = \int_{\Sigma} \int_{\Sigma} \chi(|x - y|) \sqrt{|x - y|} \partial_s u(x) \partial_s \varphi(y) d\sigma(x) d\sigma(y)$$

Justification of this form in **Collino-Joly-Lecouvez** (Waves 2013)

Cut-off function  $\chi$  used to obtain a quasi local operator.

# Comparison with Després method

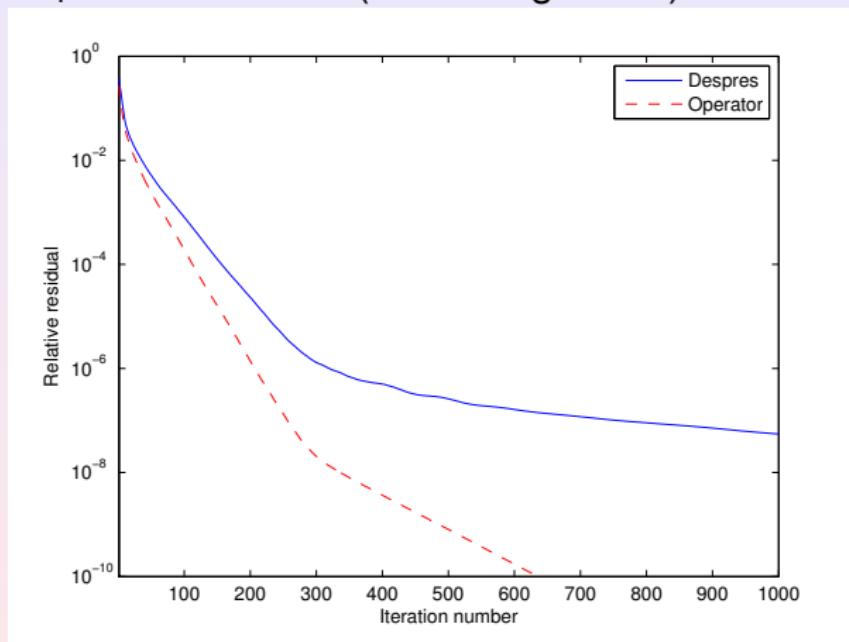
Scattering by a disc of radius  $R = 4$ , absorbing boundary condition set on  $R = 6$ ,  $\omega = 10\pi$



Subdomains are concentric discs

# Comparison with Després method

Evolution of the residual for Després method ( $z=i$ ,  $T=0$ ) and optimized method with  $\mathbb{P}_4$  finite elements (Jacobi algorithm).



⇒ Geometrical convergence for the optimized method.

# Influence of cut-off function

Number of iterations versus the radius of the cut-off function ( $\mathbb{Q}_4$ )

Radius	Jacobi	Gmres(50)
$\frac{\lambda}{8}$	497	144
$\frac{\lambda}{4}$	365	122
$\frac{\lambda}{2}$	342	114
$\lambda$	403	120
$2\lambda$	398	120
$4\lambda$	398	120

$\lambda$  is here the wavelength.

# Influence of the mesh size

Number of iterations versus the number of dofs per wavelength.

N	Jacobi	Gmres(50)
2	723	370
4	394	113
8	398	120
16	401	118
32	401	114
64	401	110

# Influence of the number of subdomains

Number of iterations versus the number of subdomains ( $\mathbb{Q}_4$ ).

Number of sub-domains	Jacobi	Gmres(50)
2	398	120
4	456	193
8	1322	435
16	3776	962

Computations performed with the same coefficients  $z, c$ . choosing different coefficients might reduce substantially the number of iterations

# Influence of the frequency

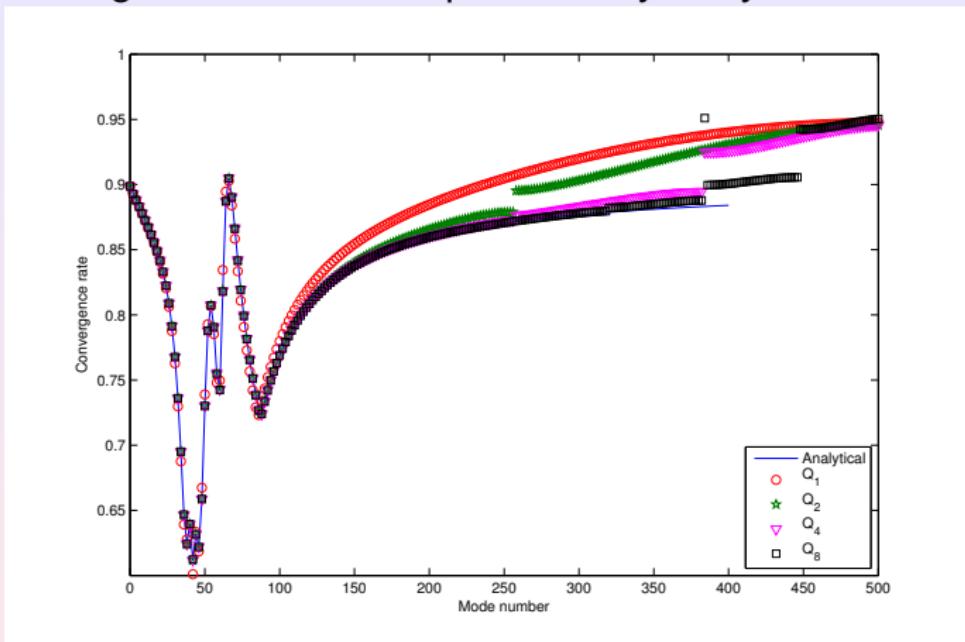
Number of iterations versus the pulsation  $\omega$  ( $\mathbb{Q}_4$  with eight degrees of freedom per wavelength)

$\omega$	Jacobi	Gmres(50)
$\pi$	150	62
$2\pi$	202	72
$4\pi$	266	86
$6\pi$	280	99
$10\pi$	398	120
$20\pi$	515	149
$40\pi$	746	189

⇒ The influence of the frequency is rather mild on this case.

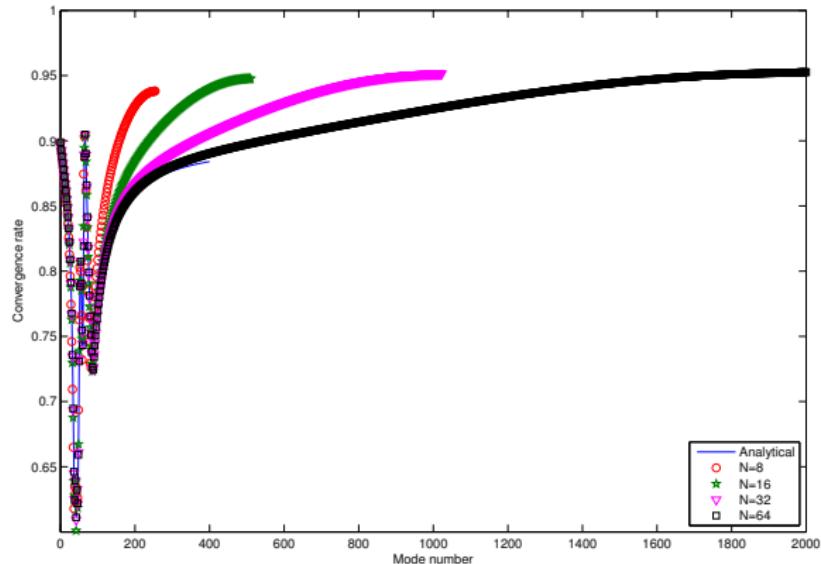
# Comparison with analytical rate

Rate of convergence can be computed analytically with modes  $e^{im\theta}$



For high values of  $m$ , the rate differ because of discretization error.  
⇒ Maximal rate of about 0.95 instead of 0.9 for the analytical computation

# Comparison with analytical rate



Final rate of 0.95 rather independant of the number of points  $N$  per wavelength.

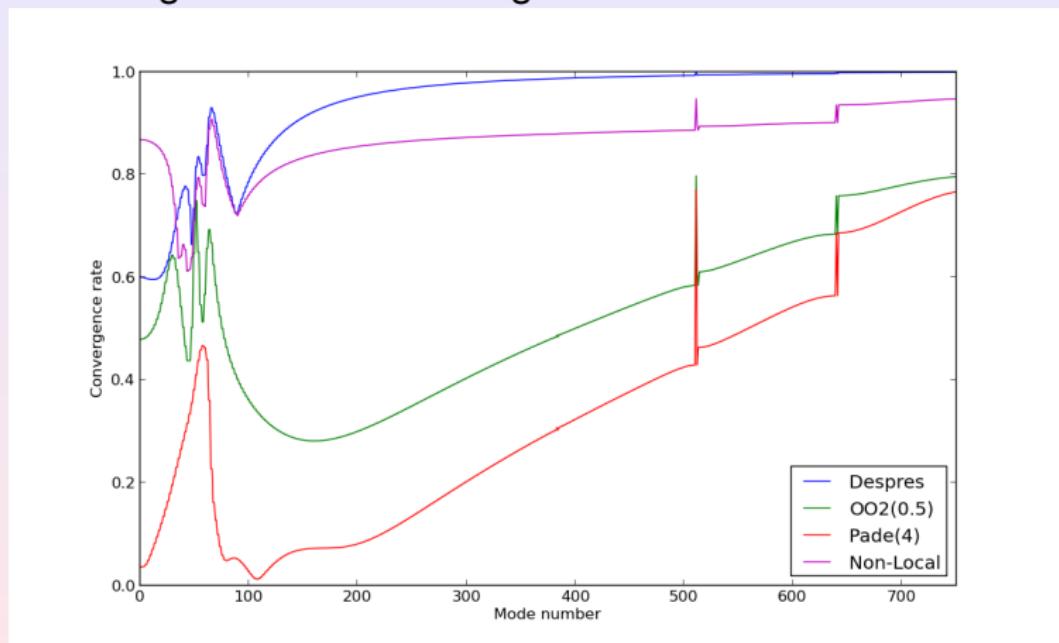
# Comparison with other transmission conditions

Comparison with the following transmission conditions

- Després operator : Després (1990)
- OO2 : Optimized second-order operator Gander, Magoules, Nataf (2002)
- Padé(N) : Square root operator approximated by Padé expansion, Antoine, Boubendir, Geuzaine (2012)
- Non-local : our approach with  $T = \Lambda\Lambda^*$

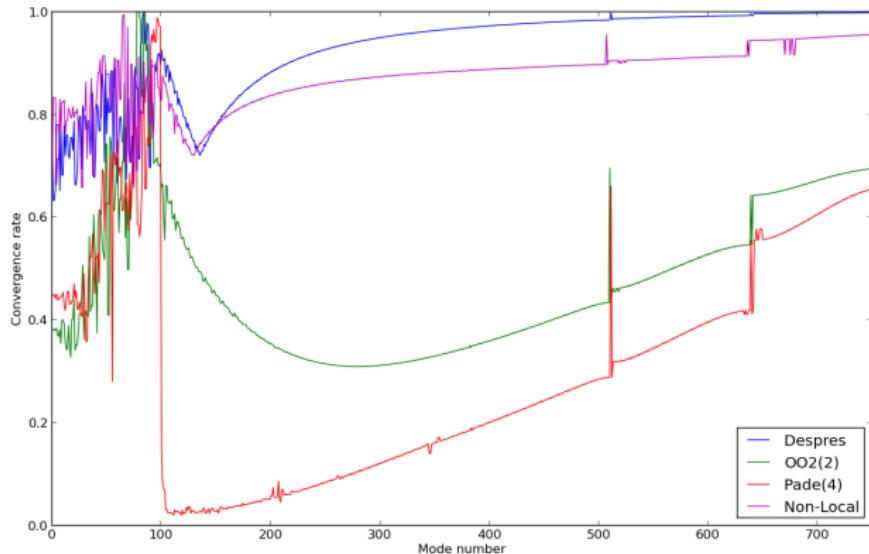
# Comparison with other transmission conditions

Rate of convergence for the homogeneous disk



# Comparison with other transmission conditions

Rate of convergence for a dielectric square ( $\rho = 2.25$ ,  $\mu = 1$ )



# Comparison with other transmission conditions

Number of iterations with Gmres(50) for the dielectric square to reach a relative residual lower than  $10^{-8}$  ( $\mathbb{Q}_6$ )

Frequency	0.01	0.1	1.0	4.0
Després	128	48	38	247
OO2	26	29	40	200
Padé(4)	94	18	34	160
Non-local	28	33	46	194

- Adjonction of local operators in  $T$
- Implement and test 3-D cases
- Case of intersecting interfaces
- Extension to 3-D Maxwell's equations