# **Optimal rotary control of the cylinder wake using POD reduced order model**

Michel Bergmann, Laurent Cordier & Jean-Pierre Brancher

Michel.Bergmann@ensem.inpl-nancy.fr

Laboratoire d'Énergétique et de Mécanique Théorique et Appliquée UMR 7563 (CNRS - INPL - UHP) ENSEM - 2, avenue de la Forêt de Haye BP 160 - 54504 Vandoeuvre Cedex, France





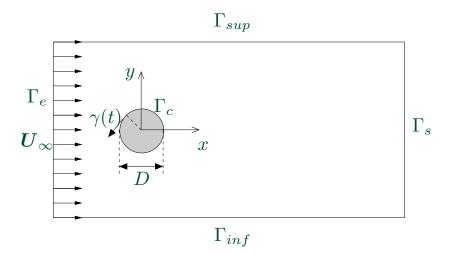
- I Flow configuration and numerical methods
- II Optimal control
- III Proper Orthogonal Decomposition (POD)
- IV Reduced Order Model of the cylinder wake (ROM)
- V Optimal control formulation applied to the ROM
- VI Results of POD ROM
- **VII Discussion**
- VIII Nelder-Mead Simplex method
- **Conclusions and perspectives**





# I - Configuration and numerical method

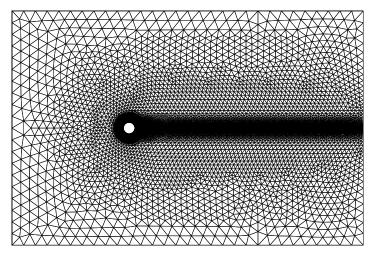
- Two dimensional flow around a circular cylinder at  $R_e = 200$
- Viscous, incompressible and Newtonian fluid
- Cylinder oscillation with a tangential velocity  $\gamma(t)$



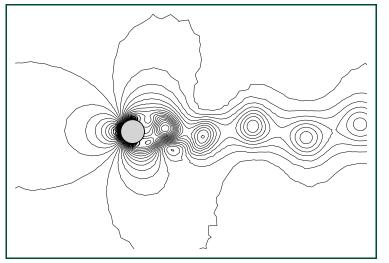
Fractional step method in time
 Finite Element Method (FEM) in space (P<sub>1</sub>)



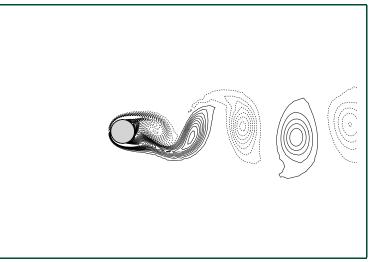
Numerical code written by M.Braza (IMFT-EMT2) & D.Ruiz (ENSEEIHT)



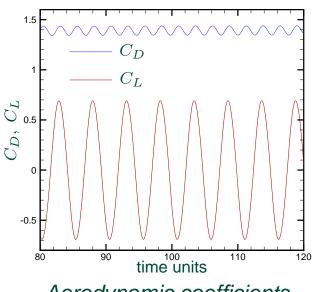
## I - Configuration and numerical method



Iso pressure at t = 100.



Iso vorticity at t = 100.



Aerodynamic coefficients.

Authors	$S_t$	$C_D$
Braza <i>et al.</i> (1986)	0.2000	1.4000
Henderson <i>et al.</i> (1997)	0.1971	1.3412
He <i>et al.</i> (2000)	0.1978	1.3560
this study	0.1983	1.3972

Strouhal number and drag coefficient.



Mathematical method allowing to determine without a priori knowledge a control law based on the optimization of a cost functional.

State equations  $\mathcal{F}(\phi, c) = 0$ ; (Navier-Stokes + I.C. + B.C.)

- Control variables c;
   (Blowing/suction, design parameters ...)
- Cost functional  $\mathcal{J}(\phi, c)$ . (Drag, lift, target function, ...)



Find a control law c and state variables  $\phi$  such that the cost functional  $\mathcal{J}(\phi, c)$  reach an extremum under the constraint  $\mathcal{F}(\phi, c) = 0$ .



# **II - Optimal control** Lagrange multipliers

Constrained optimization  $\Rightarrow$  unconstrained optimization

- ► Introduction of Lagrange multipliers  $\xi$  (adjoint variables).
- ► Lagrange functional :

$$\mathcal{L}(\phi, c, \xi) = \mathcal{J}(\phi, c) - \langle \mathcal{F}(\phi, c), \xi \rangle$$

Force  $\mathcal{L}$  to be stationary  $\Rightarrow$  look for  $\delta \mathcal{L} = 0$ :

$$\delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial c} \delta c + \frac{\partial \mathcal{L}}{\partial \xi} \delta \xi = 0$$

► Hypothesis :  $\phi$ , c and  $\xi$  assumed to be independent of each other :



$$\frac{\partial \mathcal{L}}{\partial \phi} \delta \phi = \frac{\partial \mathcal{L}}{\partial c} \delta c = \frac{\partial \mathcal{L}}{\partial \xi} \delta \xi = 0$$

# **II - Optimal control** *Optimality system*

► State equations 
$$\left(\frac{\partial \mathcal{L}}{\partial \xi}\delta\xi = 0\right)$$
:  
 $\mathcal{F}(\phi, c) = 0$   
► Co-state (adjoint) equations  $\left(\frac{\partial \mathcal{L}}{\partial \phi}\delta\phi = 0\right)$ :  
 $\left(\frac{\partial \mathcal{F}}{\partial \phi}\right)^* \xi = \left(\frac{\partial \mathcal{J}}{\partial \phi}\right)^*$   
► Optimality condition  $\left(\frac{\partial \mathcal{L}}{\partial c}\delta c = 0\right)$ :  
 $\left(\frac{\partial \mathcal{J}}{\partial c}\right)^* = \left(\frac{\partial \mathcal{F}}{\partial c}\right)^* \xi$ 



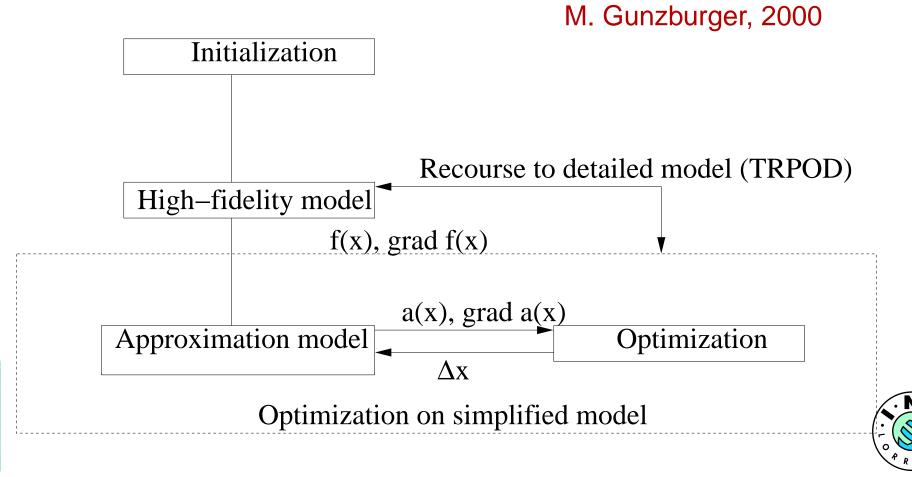
 $\Rightarrow$  Expensive method in CPU time and storage memory for large system!

⇒ Ensure only a local (generally not global) minimum



## **II - Optimal control** Reduced Order Model (ROM)

"without an inexpensive method for reducing the cost of flow computation, it is unlikely that the solution of optimization problems involving the three dimensional unsteady Navier-Stokes system will become routine"



## **II - Proper Orthogonal Decomposition (POD)**

▶ Introduced in fluid mechanics (turbulence context) by Lumley (1967).

► Look for a realization  $\phi(X)$  which is closer, in an average sense, to the realizations u(X).  $(X = (x, t) \in D = \Omega \times \mathbb{R}^+)$ 

 $\phi(X) \text{ solution of the problem :} \qquad \max_{\phi} \langle |(u, \phi)|^2 \rangle \quad \text{s.t.} \quad \|\phi\|^2 = 1.$ 

► Snapshots method, Sirovich (1987) :

$$\int_{T} C(t, t') a^{(n)}(t') dt' = \lambda^{(n)} a^{(n)}(t).$$

- ► Optimal convergence  $L^2$  norm (energy) of  $\phi(\mathbf{X})$ ⇒ Dynamical order reduction is possible.
- Decomposition of the velocity field :



$$u(x,t) = \sum_{i=1}^{N_{POD}} a^{(i)}(t) \phi^{(i)}(x).$$



#### **III - Reduced Order Model of the cylinder wake (ROM)**

► Galerkin projection of *NSE* on the POD basis :

$$\left(\boldsymbol{\phi}^{(i)}, \, \frac{\partial \boldsymbol{u}}{\partial t} + (\boldsymbol{u} \cdot \boldsymbol{\nabla})\boldsymbol{u}\right) = \left(\boldsymbol{\phi}^{(i)}, \, -\boldsymbol{\nabla}p + \frac{1}{Re}\Delta\boldsymbol{u}\right).$$

► Integration by parts (Green's formula) leads :

$$\left( \boldsymbol{\phi}^{(i)}, \, \frac{\partial \boldsymbol{u}}{\partial t} + (\boldsymbol{u} \cdot \boldsymbol{\nabla}) \boldsymbol{u} \right) = \left( p, \, \boldsymbol{\nabla} \cdot \boldsymbol{\phi}^{(i)} \right) - \frac{1}{Re} \left( (\boldsymbol{\nabla} \otimes \boldsymbol{\phi}^{(i)})^T, \, \boldsymbol{\nabla} \otimes \boldsymbol{u} \right) \\ - \left[ p \, \boldsymbol{\phi}^{(i)} \right] + \frac{1}{Re} \left[ (\boldsymbol{\nabla} \otimes \boldsymbol{u}) \boldsymbol{\phi}^{(i)} \right].$$

with 
$$[a] = \int_{\Gamma} a \cdot n \, d\Gamma$$
 and  $(A, B) = \int_{\Omega} A : B \, d\Omega = \sum_{i, j} \int_{\Omega} A_{ij} B_{ji} \, d\Omega$ .



#### **III - Reduced Order Model of the cylinder wake (ROM)**

• Velocity decomposition with  $N_{POD}$  modes :

$$\boldsymbol{u}(\boldsymbol{x},t) = \boldsymbol{u}_m(\boldsymbol{x}) + \gamma(t) \, \boldsymbol{u}_c(\boldsymbol{x}) + \sum_{k=1}^{N_{POD}} a^{(k)}(t) \boldsymbol{\phi}^{(k)}(\boldsymbol{x}).$$

▶ Reduced order dynamical system where only  $N_{gal}$  ( $\ll N_{POD}$ ) modes are retained (state equations) :

$$\frac{d a^{(i)}(t)}{d t} = \mathcal{A}_{i} + \sum_{j=1}^{N_{gal}} \mathcal{B}_{ij} a^{(j)}(t) + \sum_{j=1}^{N_{gal}} \sum_{k=1}^{N_{gal}} \mathcal{C}_{ijk} a^{(j)}(t) a^{(k)}(t) + \mathcal{D}_{i} \frac{d \gamma}{d t} + \left( \mathcal{E}_{i} + \sum_{j=1}^{N_{gal}} \mathcal{F}_{ij} a^{(j)}(t) \right) \gamma + \mathcal{G}_{i} \gamma^{2}$$
$$a^{(i)}(0) = (\boldsymbol{u}(\boldsymbol{x}, 0), \boldsymbol{\phi}^{(i)}(\boldsymbol{x})).$$

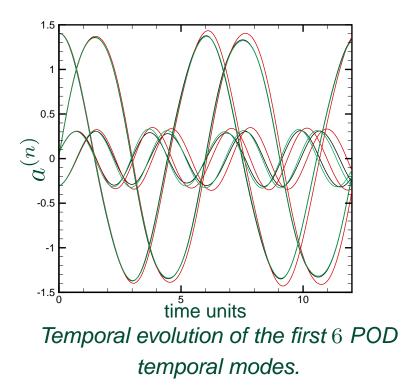


 $\mathcal{A}_i, \mathcal{B}_{ij}, \mathcal{C}_{ijk}, \mathcal{D}_i, \mathcal{E}_i, \mathcal{F}_{ij} \text{ and } \mathcal{G}_i \text{ depend on } \phi, u_m, u_c \text{ and } Re.$ 



## **IV - Reduced Order Model of the cylinder wake** *Stabilization*

Integration and "optimal" stabilization of the POD ROM for  $\gamma = A \sin(2\pi S_t t)$ , A = 2 and  $S_t = 0.5$ . POD reconstruction errors  $\Rightarrow$  temporal modes amplification



► Causes :

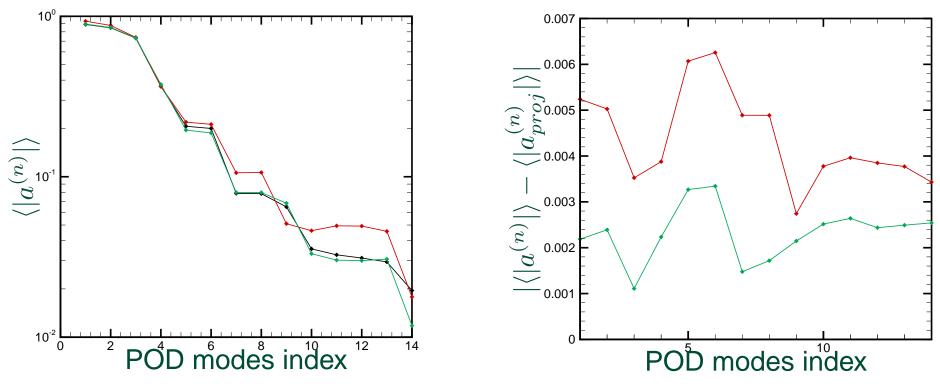
- Extraction by POD only of the large energetic eddies
- Dissipation takes place in small eddies
- ► Solution :
- Addition of an optimal artificial viscosity on each POD mode



projection (Navier-Stokes) prediction before stabilization (POD ROM) prediction after stabilization (POD ROM).



## **IV - Reduced Order Model of the cylinder wake** *Stabilization*



Comparison of energetic spectrum.

Comparison of absolute errors.

Good agreements between POD ROM spectrum and DNS spectrum



 $\Rightarrow$  Validation of the POD ROM



## **V** - Optimal control formulation based on ROM

► Objective functional :

$$\mathcal{J}(\boldsymbol{a},\gamma(t)) = \int_0^T J(\boldsymbol{a},\gamma(t)) \, dt = \int_0^T \left(\sum_{i=1}^{N_{gal}} a^{(i)^2} + \frac{\alpha}{2}\gamma(t)^2\right) \, dt.$$

 $\alpha$  : regularization parameter (penalization).

► Co-state equations :

$$\begin{cases} \frac{d\xi^{(i)}(t)}{dt} = -\sum_{j=1}^{N_{gal}} \left( \mathcal{B}_{ji} + \gamma \,\mathcal{F}_{ji} + \sum_{k=1}^{N_{gal}} \left( \mathcal{C}_{jik} + \mathcal{C}_{jki} \right) a^{(k)} \right) \xi^{(j)}(t) - 2a^{(i)} \\ \xi^{(i)}(T) = 0. \end{cases}$$

Optimality condition (gradient) :



$$\delta\gamma(t) = -\sum_{i=1}^{N_{gal}} \mathcal{D}_i \frac{d\xi^{(i)}}{dt} + \sum_{i=1}^{N_{gal}} \left( \mathcal{E}_i + \sum_{j=1}^{N_{gal}} \mathcal{F}_{ij} a^{(j)} + 2\mathcal{G}_i \gamma(t) \right) \xi^{(i)} + \alpha\gamma.$$

T R A

# **V** - Optimal control formulation based on ROM

►  $\gamma^{(0)}(t)$  done; for n = 0, 1, 2, ... and while a convergence criterium is not satisfied, do :

1. From t = 0 to t = T solve the state equations with  $\gamma^{(n)}(t)$ ;  $\hookrightarrow$  state variables  $a^{(n)}(t)$ 

2. From t = T to t = 0 solve the co-state equations with  $a^{(n)}(t)$ ;  $\hookrightarrow$  *co-state variables*  $\xi^{(n)}(t)$ 

- 3. Solve the optimality condition with  $a^{(n)}(t)$  and  $\xi^{(n)}(t)$ ;  $\hookrightarrow$  objective gradient  $\delta\gamma^{(n)}(t)$
- 4. New control law  $\hookrightarrow \gamma^{(n+1)}(t) = \gamma^{(n)}(t) + \omega^{(n)} \, \delta \gamma^{(n)}(t)$







▶ No reactualization of the POD basis.

► The energetic representativity is *a priori* different to the dynamical one :

 $\hookrightarrow$  possible inconvenient for control,

 $\hookrightarrow$  a POD dynamical system represents *a priori* only the dynamics (and its vicinity) used to build the low dynamical model.

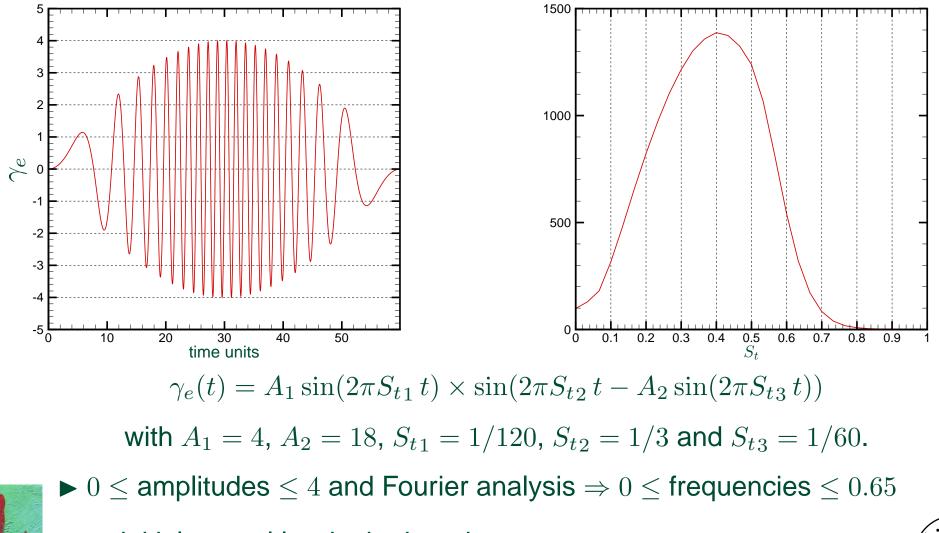
Construction of a POD basis representative of a large range of dynamics :

 $\hookrightarrow$  excitation of a great number of degrees of freedom scanning  $\gamma(t)$  in amplitudes and frequencies.





# **VI - Closed loop results** *Excitation*

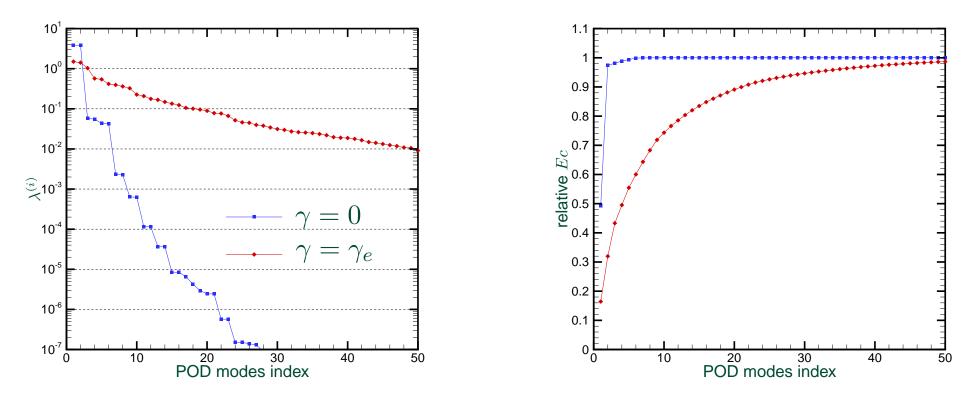




 $\blacktriangleright \gamma_e$  initial control law in the iterative process.



# **VI - Closed loop results** *Energy*



Stationary cylinder  $\gamma = 0 : \hookrightarrow 2$  modes out of 100 are sufficient to restore 97% of the kinetic energy.

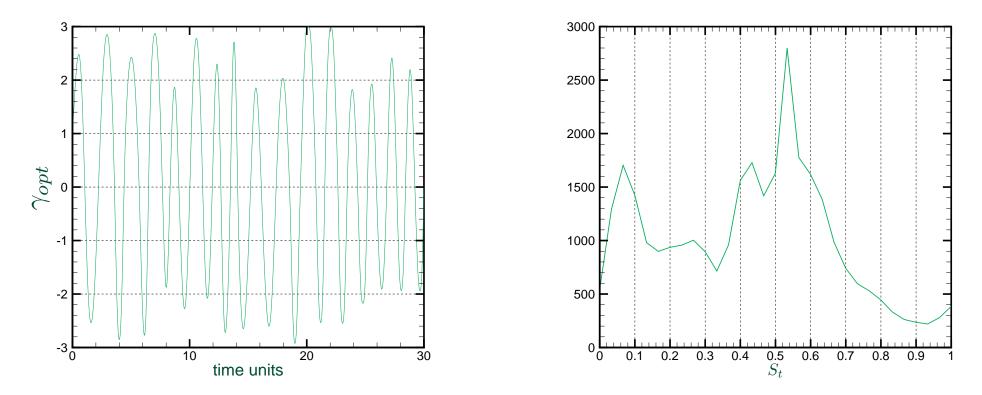


► Controlled cylinder  $\gamma = \gamma_e : \hookrightarrow 40$  modes out of 100 are then necessary to restore 97% of the kinetic energy

 $\Rightarrow$  Improvement of the POD ROM robustness to dynamical evolutions



# **VI - Closed loop results** *Optimal control*



► Reduction of the wake instationarity.  $\gamma_{opt} \simeq A \sin(2\pi S_t t)$  with A = 2.2and  $S_t = 0.53$ 

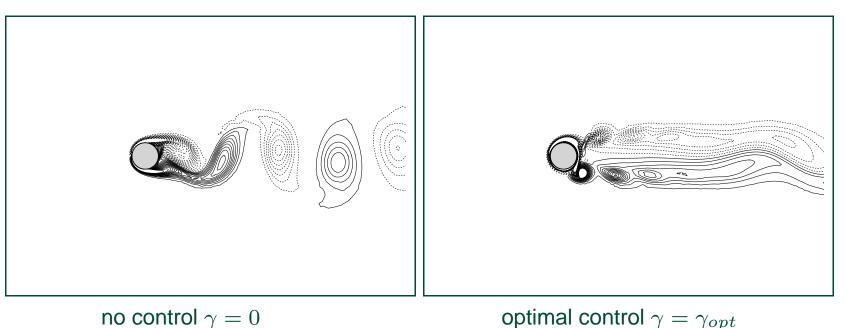
$$\mathcal{J}(\gamma_e) = 9.81 \implies \mathcal{J}(\gamma_{opt}) = 5.63.$$



The control is optimal for the reduced order model based on POD.
 Is it also optimal for the Navier-Stokes model?

# **VI - Closed loop results** Comparison of wakes' organization

#### ► No mathematical proof concerning the Navier Stokes optimality.



Isocontours of vorticity  $\omega_z$ .

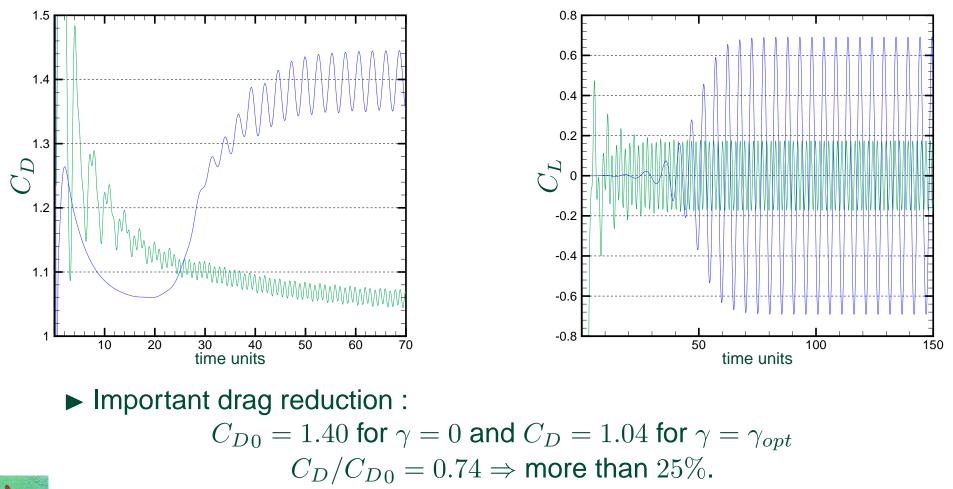
- ▶ no control :  $\gamma = 0 \Rightarrow$  Asymmetric flow.
  - $\hookrightarrow$  Large and energetic eddies.



- optimal control :  $\gamma = \gamma_{opt} \Rightarrow$  Symmetrization of the (near) wake.
  - $\hookrightarrow$  Smaller and lower energetic eddies.



# **VI - Closed loop results** Aerodynamic coefficients





► Decrease of the lift amplitude :

 $C_L = 0.68$  for  $\gamma = 0$  and  $C_L = 0.13$  for  $\gamma = \gamma_{opt}$ .



# **VI - Closed loop results** *Numerical costs*

▶ Optimal control of NSE by He *et al.* (2000) :

 → harmonic control law with A = 3 and S<sub>t</sub> = 0.75.
 ⇒ 30% drag reduction.

 ▶ Optimal control POD ROM (this study) :

 → harmonic control law with A = 2.2 and S<sub>t</sub> = 0.53.

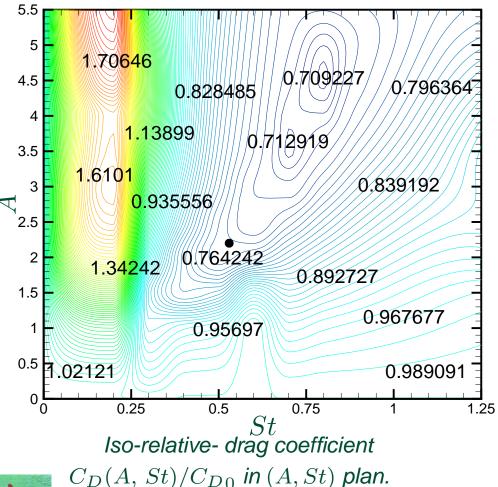
 $\Rightarrow 25\%$  drag reduction.

- Less energetic costs (greater energetic gain ?)
- Reduction costs using POD ROM compared to NSE :
  - calculus time : 100
  - Memory storage : 600



 $\hookrightarrow$  "Optimal" control of 3D flows becomes possible !





#### **Observations**

Minimum is located in a smooth valley

 $\hookrightarrow$  Global minimum : around A=4.4 and St=0.76

 Maximum is located in a sharp peak

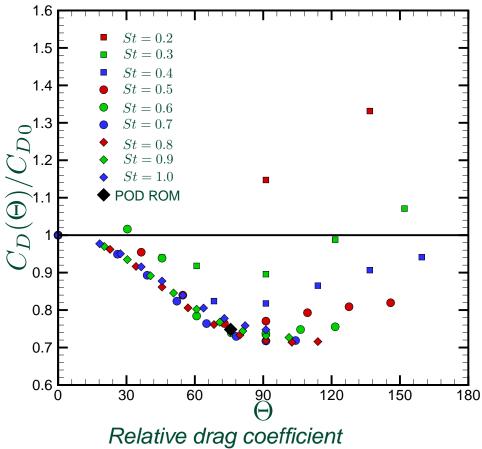
 $\hookrightarrow$  Global maximum : near St=0.2, the natural frequency : lock-on flow



Finding the global minimum with an optimization algorithm may be difficult due to the smooth valley



# VII - Discussion Maximum angle of rotation



vs. maximum angle of rotation.

• Maximum angle of rotation :  $\Theta = \max_t \left\{ \theta(t) \right\} = \frac{A}{\pi St}$ 

#### Observations

No drag reduction possible near natural frequency

► Maximum drag reduction around  $\Theta_{max} = 95^{\circ}$ 

 $\hookrightarrow$  For all frequencies g.t. natural frequency

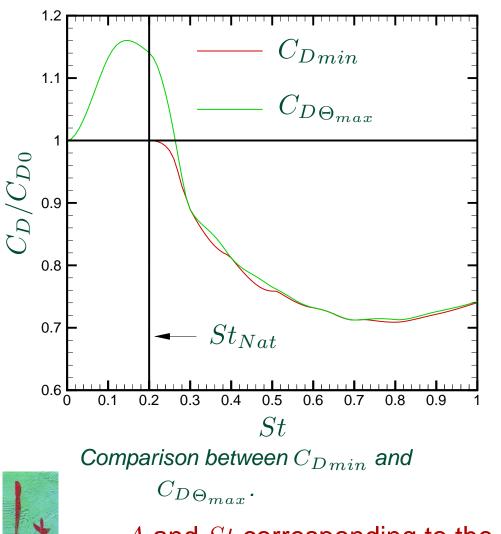
 $\hookrightarrow$  Minimum drag :

$$C_D = 0.71 \times C_{D0} = 0.98$$



Existence of an "optimal" maximum angle of rotation  $\Theta_{max}$ .





#### **Notations**

$$C_{D\min}(St) = \min_{A \in \mathbb{R}} C_D(\Theta, St)$$
$$C_{D\Theta_{\max}}(St) = C_D(\Theta_{\max}, St)$$

#### **Observations**

Good agreements between  $C_{Dmin}$  and  $C_{D\Theta_{max}}$  for  $St > St_{Nat}$ 

 $\blacktriangleright \Theta_{max}$  is not optimal for St <  $St_{Nat}$ 



A and St corresponding to the minimal drag seems dependent :  $A/St = 5.2 \ (\Theta_{max} = 95^{\circ}).$ 



POD ROM control law does not correspond to the global minimum

 $\hookrightarrow$  POD ROM parameters : A = 2.2 and St = 0.53 ( $\Theta = 76^{\circ}$ )  $\Rightarrow C_D = 1.04$ 

 $\hookrightarrow \text{Global minimum parameters} : A = 4.4 \text{ and } St = 0.76$ ( $\Theta = 105^{\circ} \neq \Theta_{max} = 95^{\circ}$ )  $\Rightarrow C_D = 0.98$ 

▶ Results in (A, St) quite different but not so far in terms of  $C_D$ 

 $\hookrightarrow$  The smooth valley is reached

Improvement : coupling to the POD ROM approach an efficient new optimization algorithm for smooth fonctions



 $\hookrightarrow$  Take results obtained by POD ROM as initial conditions

#### Advantages

- Numerical simplicities
- Adaptive topology
- Gradients calculations not necessary
- Good results with smooth functions

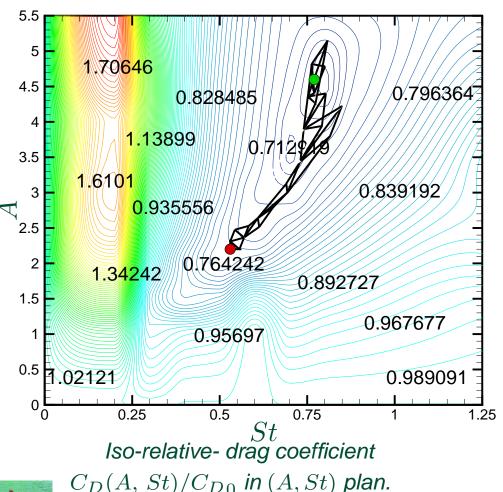
#### Drawbacks

- ► No proof of optimality for simplex dimensions greater than two
- Need to fix free parameters



Maybe more iterations than gradient based optimisation algorithms...





#### **Observations**

Topology adaptation function of the curve of the valley

► Minimum found by Nelder-Mead simplex method : A = 4.5 and  $St = 0.76 \Rightarrow \Theta = 108^{\circ}$  $\hookrightarrow$  Seems to be the global minimum

▶ 30 NSE resolutions  $\Rightarrow 5\%$  additive drag reduction compared to POD ROM



Relative drag reduction by POD ROM : 25% (1 NSE resolution) Usefulness of coupling a new algorithm?



#### Conclusions

- Important drag reduction obtained by POD ROM : more than 25% of relative drag reduction
- This solution is not the global minimum of the drag function
- POD ROM compared to NSE ⇒ important reduction of numerical costs :
  - $\hookrightarrow$  Reduction factor of the calculus : 100
  - $\hookrightarrow$  Reduction factor of the memory storage : 600

"OPTIMAL" CONTROL OF 3D FLOWS POSSIBLE BY POD ROM

- Existence of an optimal maximum angle of rotation for effective drag reduction,  $\Theta_{max}=95^\circ$
- Coupling POD ROM with the Nelder-Mead simplex method leads a priori to the global minimum of the drag function



 But the gain on the drag function is quite small compared to result obtained by POD ROM



#### **Perspectives**

Improve the representativity of the POD ROM

 $\hookrightarrow$  "Optimize" the temporal excitation  $\gamma_e$  $\hookrightarrow$  Mix snapshots corresponding to different dynamics (temporal excitations)

- Look for harmonic control  $\gamma(t) = A \sin(2\pi S_t t)$  with POD basis reactualization (closed loop on NSE and not only on POD ROM)
- Coupling the POD ROM approach with Trust Region Methods (TRPOD)

 $\implies$  proof of convergence under weak conditions

- Introducing the pressure into the POD dynamical system
  - $\hookrightarrow$  pressure contribution to drag coefficient : 80%



Optimal control of the Navier-Stokes equations

