

# A dynamical approach of some Hamilton-Jacobi equations using weak KAM theory

Philippe Thieullen (Bordeaux)

(Joint work with Xifeng Su, BNU Beijing)

**April 2018**

**School of Mathematics and Statistics  
University of New South Wales**

## Summary of the talk

- I. A summary of weak KAM theory
- II. Discrete weak KAM approach
- III. Numerical analysis

# I. Summary

# I. Summary - Main goal

## Apply some methods in dynamical systems to solve

- $\frac{\partial u}{\partial t}(t, x) + H(x, \nabla u(t, x)) = 0$ , (non stationary Hamilton-Jacobi),
- $H(x, \nabla \bar{u}(x)) = \bar{H}$ , (cell equation),
- $\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0$ , (discounted cell equation)

where  $H(x, p) : \mathbb{T}^d \times \mathbb{R}^d \rightarrow \mathbb{R}$  is  $C^2$ , periodic in  $x$ , strictly convex and super-linear in  $p$ , autonomous

- $\left[ \frac{\partial^2 H(x, p)}{\partial p_i \partial p_j} \right] \geq \alpha [\delta_{ij}]$ ,
- $\lim_{R \rightarrow +\infty} \inf_{x, \|p\| \geq R} \frac{H(x, p)}{\|p\|} = +\infty$

$H(x, p)$  is called a *Tonelli Hamiltonian*

# I. Summary - known results

## Solutions in the viscosity sense

Consider the cell equation

$$H(x, \nabla \bar{u}(x)) = \bar{H}$$

$u$  is said to be a sub-solution in the viscosity sense if

- $\forall x_0 \in \mathbb{T}^d, \forall \phi \in C^1(\mathbb{T}^d, \mathbb{R})$ , if  $\phi(x_0) = u(x_0)$  and  $\forall x, \phi(x) \geq u(x)$ ,
- then  $H(x_0, \nabla \phi(x_0)) \leq \bar{H}$ .

**Existence of solutions** In all 3 equations a solution exists in the viscosity sense

- $\frac{\partial u}{\partial t}(t, x) + H(x, \nabla u(t, x)) = 0,$   $u(t, x)$  is unique,
- $H(x, \nabla \bar{u}(x)) = \bar{H},$   $\bar{u}(x)$  is not unique,  $\bar{H}$  is unique,
- $\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0,$   $u_\delta(x)$  is unique.

# I. Summary - Solutions

## Explicit solution of the non stationary Hamilton-Jacobi equation

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) + H(x, \nabla u(t, x)) = 0 \\ u(0, x) = u_0(x) \end{cases}$$

the unique solution is given by

$$u(t, x) = \inf_{\substack{\gamma \in C^2([-t, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \left[ u_0(\gamma(-t)) + \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right],$$

where  $L(x, v)$  is the Legendre transform of  $H(x, p)$

$$L(x, v) = \sup_p \{v \cdot p - H(x, p)\}$$

# I. Summary - Solutions

## Explicit solution of the effective energy

$$H(x, \nabla \bar{u}(x)) = \bar{H}$$

then the effective energy is given by

$$-\bar{H} := \lim_{t \rightarrow +\infty} \inf_{\gamma \in C^2([-t, 0], \mathbb{R}^d)} \left[ \frac{1}{t} \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right]$$

## Weak KAM theorem [Fathi, 1997]

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) + H(x, \nabla u(t, x)) = 0 \\ u(0, x) = u_0(x) \end{cases}$$

$$u(t, x) + t\bar{H} \rightarrow \bar{u}(x) \quad \text{in the } C^0 \text{ topology}$$

$\bar{u}(x)$  is called weak KAM solution

# I. Summary - Solutions

## Explicit solution for the discounted cell equation

$$\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0$$

the unique solution is given by

$$u_\delta(x) = \inf_{\substack{\gamma \in C^2((-\infty, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \int_{-\infty}^0 e^{s\delta} L(\gamma(s), \dot{\gamma}(s)) ds,$$

## Weak KAM theorem [Davini, Fathi, Iturriaga, Zavidovique, 2016]

$$\lim_{\delta \rightarrow 0} \left( u_\delta(x) + \frac{\bar{H}}{\delta} \right) = \bar{u}(x) \quad \text{in the } C^0 \text{ topology}$$

$$H(x, \bar{u}(x)) = \bar{H}$$



# I. Summary - Characteristics method

**Weierstrass-Tonelli solution** Assume  $u_0 \in C^2(\mathbb{T}^d)$

$$\begin{cases} \frac{\partial u}{\partial t} + H(x, \nabla u(t, x)) = 0, \\ u(0, x) = u_0(x) \end{cases}$$

Then a  $C^2$  solution exists uniformly in  $x$  for a short time  $0 \leq t \leq t_0$

$$u(t, x) := u_0(X_{t,x}(0)) + \int_0^t L(X_{t,x}(s), \dot{X}_{t,x}(s)) ds$$

where  $X_{t,x}(s)$ ,  $P_{t,x}(s)$  are solutions of the Hamiltonian flow

$$\begin{aligned} \dot{X}_{x,t} &= \frac{\partial H}{\partial p}(X_{x,t}, P_{x,t}), & X_{x,t}(t) &= x \\ \dot{P}_{x,t} &= -\frac{\partial H}{\partial x}(X_{x,t}, P_{x,t}), & P_{x,t}(0) &= \nabla u_0(X_{x,t}(0)) \end{aligned}$$

# I. Summary - Characteristics method

**Remark** The Weierstrass-Tonelli solution is a minimizer

$$u(t, x) = \inf_{\substack{\gamma \in C^2([-t, 0], \mathbb{R}^d) \\ \gamma(t) = x}} \left\{ u_0(\gamma(-t)) + \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right\}$$

**Proof**

$$u(t, \gamma(t)) - u_0(\gamma(0)) = \int_0^t \frac{d}{ds} u(s, \gamma(s)) ds = \int_0^t \left[ \frac{\partial u}{\partial t} + \nabla u \cdot \dot{\gamma} \right] ds$$

$$p \cdot v \leq H(x, p) + L(x, v),$$

$$u(t, \gamma(t)) - u_0(\gamma(0)) \leq \int_0^t \left[ \frac{\partial u}{\partial t} + H(\gamma, \nabla u) + L(\gamma, \dot{\gamma}) \right] ds = \int_0^t L(\gamma, \dot{\gamma}) ds$$

# I. Summary - Dynamic programming

## Dynamic programming for the discounted equation

$$\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0$$

$$u_\delta(x) = \inf_{\substack{\gamma \in C^2((-\infty, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \int_{-\infty}^0 e^{s\delta} L(\gamma(s), \dot{\gamma}(s)) ds$$

$$u_\delta(x) = \inf_{\substack{\gamma \in C^2([-t, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \left\{ e^{-t\delta} u_\delta(\gamma(-t)) + \int_{-t}^0 e^{s\delta} L(\gamma(s), \dot{\gamma}(s)) ds \right\}$$

## Dynamic programming for the cell equation

$$u_\delta(x) + \frac{\bar{H}}{\delta} \rightarrow \bar{u}(x), \quad H(\bar{u}(x), \nabla \bar{u}(x)) = \bar{H}$$

$$\bar{u}(x) - t\bar{H} = \inf_{\substack{\gamma \in C^2([-t, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \left\{ \bar{u}(\gamma(-t)) + \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right\}$$

# I. Summary - Lax-Oleinik operator

**Definition - Lax-Oleinik operator** For every  $u_0 \in C^0(\mathbb{T}^d)$

- $T^t[u](x) := \inf_{\substack{\gamma \in C^2([-t,0], \mathbb{R}^d) \\ \gamma(0)=x}} \left\{ u(\gamma(-t)) + \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right\}$
- $T^{s+t} = T^s \circ T^t$

**Weak KAM Theorem [Fathi]** for every  $u \in C^0(\mathbb{T}^d)$

- $u(t, x) := T^t[u_0](x)$  solves 
$$\begin{cases} \frac{\partial u}{\partial t} + H(x, \nabla u(t, x)) = 0, \\ u(0, x) = u_0(x) \end{cases}$$
- $T^t[u_0] + t\bar{H} \rightarrow \bar{u}(x)$  (in the  $C^0$  topology)
- $H(x, \nabla \bar{u}(x)) = \bar{H} \iff T^t[\bar{u}] = \bar{u}(x) - t\bar{H}, \quad \forall t \geq 0$

## II. Discrete weak KAM approach

## II. The discrete Lax-Oleinik operator

### The cell equation

$$H(x, \nabla \bar{u}(x)) = \bar{H}, \quad H(x, p) \text{ is a Tonelli Hamiltonian}$$

**The (continuous) Lax-Oleinik operator** For every  $u \in C^0(\mathbb{T}^d)$

$$T^t[u](x) := \inf_{\substack{\gamma \in C^2([-t, 0], \mathbb{R}^d) \\ \gamma(0) = x}} \left\{ u(\gamma(-t)) + \int_{-t}^0 L(\gamma, \dot{\gamma}) ds \right\}$$

**The discrete Lax-Oleinik operator**

$$E_\tau(x, y) := \tau L\left(x, \frac{y - x}{\tau}\right), \quad \tau > 0 \quad (\text{small})$$

$$T_\tau[u](y) = \min_{x \in \mathbb{R}^d} \left\{ u(x) + E_\tau(x, y) \right\}, \quad \forall u \in C^0(\mathbb{T}^d)$$

## II. The discrete cell equation

### Notations

- $H(x, p)$  a Tonelli Hamiltonian
- $L(x, v) = \sup_{p \in \mathbb{R}^d} \{p \cdot v - H(x, p)\}$
- $E_\tau(x, y) = \tau L(x, \frac{y-x}{\tau})$
- $T_\tau[u](y) = \min_x \{u(x) + E_\tau(x, y)\}, \quad \forall u \in C^0(\mathbb{T}^d)$

### Theorem (Su, Th): first part

- $\exists(\bar{u}_\tau, \bar{E}_\tau), \quad T_\tau[\bar{u}_\tau] = \bar{u}_\tau + \bar{E}_\tau \quad (\text{non unique } \bar{u}_\tau \in C^0(\mathbb{T}^d))$
- $\bar{E}_\tau = \lim_{n \rightarrow +\infty} \inf_{x_0, x_1, \dots, x_n} \frac{1}{n} \sum_{k=1}^n E_\tau(x_{k-1}, x_k) \quad (\text{unique})$
- $\bar{E}_\tau = -\tau \bar{H} + O(\tau^2) \quad (\text{as } \tau \rightarrow 0)$
- $\bar{u}_\tau \rightarrow \bar{u}, \quad (\text{for some subsequence, } \text{Lip}(\bar{u}_\tau) \leq C)$
- $H(x, \nabla \bar{u}(x)) = \bar{H} \quad (\text{in the viscosity sense})$

Main drawback: a sub-sequence need be taken for  $u$ .

## II. The discounted discrete Lax-Oleinik

### The discounted cell equation

$$\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0, \quad (\text{in the viscosity sense})$$

$$u_\delta(x) = \inf_{\gamma(0)=x} \int_{-\infty}^0 e^{s\delta} L(\gamma(s), \dot{\gamma}(s)) ds$$

### The discounted Lax-Oleinik operator

$$T_\delta^t[u](x) = \inf_{\gamma(0)=x} \left\{ e^{-t\delta} u(\gamma(-t)) + \int_{-t}^0 e^{s\delta} L(\gamma(s), \dot{\gamma}(s)) ds \right\}$$

### The discounted discrete Lax-Oleinik operator

$$T_{\tau,\delta}[u](y) = \min_x \left\{ (1 - \tau\delta)u(x) + E_\tau(x, y) \right\} \quad E_\tau(x, y) = \tau L\left(x, \frac{y-x}{\tau}\right)$$



## II. The discounted discrete cell equation

### Notations

- $\delta u_\delta(x) + H(x, \nabla u_\delta(x)) = 0$  (the discounted cell equation)
- $H(x, \nabla \bar{u}(x)) = \bar{H}$  (the cell equation)
- $T_{\tau, \delta}[u](y) = \min_x \{(1 - \tau\delta)u(x) + E_\tau(x, y)\}$
- $\bar{E}_\tau$  defined uniquely by  $T_\tau[\bar{u}_\tau] = \bar{u}_\tau + \bar{E}_\tau$  for some  $\bar{u}_\tau$

### Theorem (Su, Th): second part

- $\exists u_{\tau, \delta} \quad T_{\tau, \delta}[u_{\tau, \delta}] = u_{\tau, \delta}$  (unique)

$$u_{\tau, \delta}(x) = \inf_{\substack{(x-k)_{k \geq 0} \\ x_0 = x}} \sum_{k=0}^{\infty} (1 - \tau\delta)^k E_\tau(x_{-k-1}, x_{-k})$$

- $u_{\tau, \delta} = u_\delta + O\left(\frac{\tau}{\delta}\right)$  as  $\tau \rightarrow 0$  (for fixed  $\delta > 0$ )
- $u_\delta + \frac{\bar{H}}{\delta} \rightarrow \bar{u}$  (DFIZ theorem, no speed)
- $u_{\tau, \delta} - \frac{\bar{E}_\tau}{\tau\delta} \rightarrow \bar{u}_\tau$  as  $\delta \rightarrow 0$  (for fixed  $\tau > 0$ , no speed)

$$\left[ u_{\tau, \delta} - \frac{\bar{E}_\tau}{\tau\delta} \right] - \left[ u_\delta + \frac{\bar{H}}{\delta} \right] = O\left(\frac{\tau}{\delta}\right)$$

## II. Selection principle

**Recall**  $u_{\tau,\delta} - \frac{\bar{E}_\tau}{\tau\delta} \rightarrow \bar{u}_\tau$  is a distinguished solution of

$$T_\tau[\bar{u}_\tau](y) = \min_x \{\bar{u}_\tau(x) + E_\tau(x, y)\} = \bar{u}_\tau(y) + \bar{E}_\tau$$

**Stationary plan** A probability measure  $\pi(dx, dy)$  on  $\mathbb{T}^d \times \mathbb{T}^d$  s.t.

$$\iint \phi(x) \pi(dx, dy) = \iint \phi(y) \pi(dx, dy) \quad \forall \phi \in C^0(\mathbb{T}^d)$$

**Minimizing plan** A stationary plan  $\pi_{\min}$  is minimizing if

$$\bar{E}_\tau = \min_\pi \iint E_\tau(x, y) \pi(dx, dy) = \iint E_\tau(x, y) \pi_{\min}(dx, dy)$$

**Mañé Potential** Define on  $\mathbb{T}^d \times \mathbb{T}^d$

$$\Phi_\tau(x, y) = \inf_{n \geq 1} \inf_{\substack{x_0, x_1, \dots, x_n \\ x_0 = x, x_n = y}} \sum_{k=0}^{n-1} [E_\tau(x_k, x_{k+1}) - \bar{E}_\tau]$$

## II. Selection principle bis

### Notation

- $T_\tau[u_\tau](y) = \min_x \{u_\tau(x) + E_\tau(x, y)\}$
- $T_{\tau, \delta}[u](y) = \min_x \{(1 - \tau\delta)u(x) + E_\tau(x, y)\}$

**Theorem: second part** If  $T_{\tau, \delta}[u_{\tau, \delta}] = u_{\tau, \delta}$

- $u_{\tau, \delta} - \frac{\bar{E}_\tau}{\tau\delta} \rightarrow \bar{u}_\tau$  as  $\delta \rightarrow 0$  where  $T_\tau[\bar{u}_\tau] = \bar{u}_\tau + \bar{E}_\tau$
- First characterization

$$\bar{u}_\tau(z) = \sup \left\{ w(z) : T_\tau[w] = w + \bar{E}_\tau \text{ and } \iint w(y) d\pi(x, y) \leq 0, \forall \pi \right\}$$

- Second characterization

$$\bar{u}_\tau(z) = \inf \left\{ \iint \Phi_\tau(x, z) d\pi(x, y) : \pi \text{ minimizing plan} \right\}$$

## II. The discrete Aubry set

### Effective energy and Mañé potential

$$\bar{E}_\tau = \lim_{n \rightarrow +\infty} \inf_{x_0, x_1, \dots, x_n} \frac{1}{n} \sum_{k=1}^n E_\tau(x_{k-1}, x_k)$$

$$\Phi_\tau(x, y) = \inf_{n \geq 1} \inf_{\substack{x_0, x_1, \dots, x_n \\ x_0 = x, x_n = y}} \sum_{k=0}^{n-1} [E_\tau(x_k, x_{k+1}) - \bar{E}_\tau]$$

**The projected Aubry set**  $\Phi_\tau(x, x) \geq 0, \quad \forall x \in \mathbb{T}^d$

$$\text{ProjectedAubry}(E_\tau) := \{x \in \mathbb{T}^d : \Phi_\tau(x, x) = 0\}$$

**The Aubry set**  $\Phi_\tau(x, y) + \Phi_\tau(y, x) \geq 0, \quad \forall (x, y) \in \mathbb{T}^d \times \mathbb{T}^d$

$$\text{Aubry}(E_\tau) := \{(x, y) \in \mathbb{T}^d \times \mathbb{T}^d : \Phi_\tau(x, y) + \Phi_\tau(y, x) = 0\}$$

## II. Minimizing configuration

**Definition** A configuration  $(x_k)_{k=-\infty}^{+\infty}$  is minimizing if

$$\sum_{k=m+1}^{m+n} [E_\tau(x_{k-1}, x_k) - \bar{E}_\tau] \leq \sum_{k=m+1}^{m+n} [E_\tau(y_{k-1}, y_k) - \bar{E}_\tau],$$

$$\forall m \in \mathbb{Z}, n \geq 0 \quad \forall (y_k)_{k=-\infty}^{+\infty} \quad y_{m+1} = x_{m+1}, y_{m+n} = x_{m+n}$$

**Existence** Let  $T_\tau[\bar{u}_\tau] = \bar{u}_\tau + \bar{E}_\tau$ ,  $y \in \mathbb{R}^d$ , and

$$x_{-n-1} \in \arg \min_{x \in \mathbb{R}^d} \{u_\tau(x) + E_\tau(x, x_{-n})\}, \quad x_0 = y$$

- $(x_{-n})_{n \geq 0}$  is a backward minimizing configuration  $\rightarrow$  Aubry( $E_\tau$ )
- the Aubry set is a set of uniqueness for  $\bar{u}_\tau$

$$\bar{u}_\tau(y) = \min_{x \in \text{ProjectedAubry}(E_\tau)} \{\bar{u}_\tau(x) + \Phi_\tau(x, y)\}$$

## III. Numerical analysis

### III. The approximation scheme

**Solve**  $H(x, \nabla \bar{u}(x)) = \bar{H}$

- compute explicitly,  $L(x, v) := \sup_{p \in \mathbb{R}^d} [p \cdot v - H(x, p)],$
  - choose a step time  $\tau > 0$  and define  $E_\tau(x, y) := \tau L\left(x, \frac{y-x}{\tau}\right)$
  - define  $T_\tau[u_\tau](y) := \min_{x \in \mathbb{T}^d} \{u_\tau(x) + E_\tau(x, y)\}$
  - solve  $T_\tau[\bar{u}_\tau] = \bar{u}_\tau + \bar{E}_\tau$  using Ishikawa's iterative method
- $$u_\tau^{(0)} = 0, \quad u_\tau^{(n+1)} := \frac{u_\tau^{(n)} + T_\tau[u_\tau^{(n)}]}{2} - \min\left(\frac{u_\tau^{(n)} + T_\tau[u_\tau^{(n)}]}{2}\right)$$
- $u_\tau^{(n)} \rightarrow \bar{u}_\tau, \quad \min(T_\tau[u_\tau^{(n)}]) \rightarrow \bar{E}_\tau$
  - $-\frac{\bar{E}_\tau}{\tau} \rightarrow \bar{H}, \quad \bar{u}_\tau \rightarrow \bar{u}$  for some sub-sequence  $\tau \rightarrow 0$
  - $H(x, \nabla \bar{u}(x)) = \bar{H},$  in the viscosity sense

## III. The inverse pendulum

### Frenkel-Kontorova model

- $H(x, p) = \frac{1}{2}(p + \lambda)^2 - KV(x), \quad V(x) = \frac{1}{(2\pi)^2}(1 - \cos 2\pi x)$
- $L(x, v) = \frac{1}{2}v^2 - \lambda v + KV(x)$
- $E_\tau(x, y) = \tau L(x, \frac{y-x}{\tau}) = \frac{1}{2\tau}(y-x)^2 - \lambda(y-x) + \tau KV(x)$
- solve  $\bar{u}_\tau(y) + \bar{E}_\tau = \min_x \{\bar{u}_\tau(x) + E_\tau(x, y)\}$
- $-\bar{E}_\tau/\tau \rightarrow \bar{H}$
- $\bar{u}_\tau \rightarrow \bar{u}$  for some subsequence
- $\frac{1}{2}(\nabla \bar{u} + \lambda)^2 - KV(x) = \bar{H}$



### III. The continuous weak KAM

#### Explicit solutions

- $\frac{1}{2}(\nabla\bar{u} + \lambda)^2 - KV(x) = \bar{H} \Leftrightarrow \nabla\bar{u} = -\lambda \pm \sqrt{2(\bar{H} + KV(x))}$
- $\lambda_* := \int_0^1 \sqrt{2KV(s)} ds = \frac{2\sqrt{K}}{\pi^2} \simeq 0.203 \sqrt{K}$
- if  $\lambda > \lambda_*$  then  $\exists! \bar{H}$  s.t.  $\lambda = \int_0^1 \sqrt{2(\bar{H} + KV(s))} ds$

$$\forall 0 \leq x \leq 1, \quad \bar{u}(x) = -\lambda x + \int_0^x \sqrt{2(\bar{H} + KV(s))} ds$$

- If  $\lambda \leq \lambda_*$  then  $\bar{H} = 0$ ,

$$\exists! x_* \in [\frac{1}{2}, 1], \quad \lambda = \int_0^{x_*} \sqrt{2KV(s)} ds - \int_{x_*}^1 \sqrt{2KV(s)} ds$$

$$\begin{cases} \bar{u}(x) = -\lambda x + \int_0^x \sqrt{2KV(s)} ds, & \text{if } 0 \leq x \leq x_* \\ \bar{u}(x) = -\lambda x + \int_0^{x_*} \sqrt{2KV(s)} ds - \int_{x_*}^x \sqrt{2KV(s)} ds, & \text{if } x_* \leq x \leq 1 \end{cases}$$

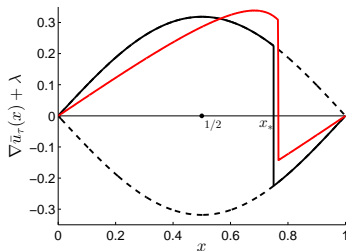
### III. The discrete weak KAM: $\bar{H} = 0$

#### The cell equation

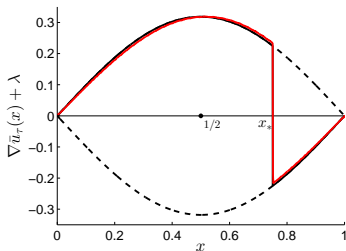
- $\frac{1}{2}(\nabla \bar{u}(x) + \lambda)^2 - KV(x) = \bar{H}$ ,  $V(x) = \frac{1}{(2\pi)^2}(1 - \cos 2\pi x)$
- $\bar{u}_\tau(y) + \bar{E}_\tau = \min_x \left\{ \bar{u}_\tau(x) + \frac{1}{2\tau}(y - x)^2 - \lambda(y - x) + \tau KV(x) \right\}$
- $K = 1$ ,  $x_* = \frac{3}{4}$ ,  $\lambda = 0.143 < \lambda_* = 0.202$ ,  $N_{\text{grid}} = 1000$

#### Numerical approximation

In black:  $\bar{u}$ , in red:  $\bar{u}_\tau$








$$\tau = 1, \bar{E}_\tau \sim 10^{-12}, N_{\text{iter}} = 32$$



$$\tau = 0.1, \bar{E}_\tau \sim 10^{-20}, N_{\text{iter}} = 196$$

# Bibliography I

-  S. Aubry and P.Y. Le Daeron, The discrete Frenkel-Kontorova model and its extensions: I. Exact results for the ground states, *Physica D*, Vol. 8 (1983), 381-422.
-  W. Chou, R.B. Griffiths, Ground states of one-dimensional systems using effective potentials *Phys. Rev. B*, Vol. 34 (1986), 6219-6234.
-  A. Fathi, The weak KAM theorem in Lagrangian dynamics, Book to appear, Cambridge University Press (see author website).
-  E. Garibaldi, Ph. Thieullen, Minimizing orbits in the discrete Aubry-Mather model. *Nonlinearity*, Vol. 24 (2011), 563-611.
-  Xifeng Su, Ph. Thieullen, Convergence of discrete Aubry-Mather model in the continuous limit, *Nonlinearity*, Vol. 31 (2018), 2126-2155.