

Zero-temperature Gibbs measures for some sub-shifts of finite type

Philippe Thieullen (Bordeaux)

Academy of Mathematics
and Systems Sciences

January 26, 2015

Summary of the talk

- I. Motivations
- II. Thermodynamic formalism
- III. Selection principle

I. Motivations

Cell equation Consider $H(x, p)$ a Tonelli Hamiltonian: C^2 , autonomous, periodic in x , super-linear in p , and definite positive. The cell equation is

$$H(x, d_x u) = \bar{H}, \quad u \text{ as a viscosity solution}$$

Objective The cell equation is very degenerate. There are two approaches:

- ♦ a PDE approach (L.P.V.),
- ♦ a dynamical approach using Fathi's weak KAM theory.
- ♦ A third approach:

\bar{H} and u are thermodynamic objects that can be obtained as a limit when the temperature of some the system goes to 0

Discretization in time (Bardi, Capuzzo-Dolcetta, Falcone, 2008)

- ◆ solve the discounted cell equation

$$\delta u_\delta + H(x, d_x u_\delta) = 0$$

- ◆ use the representation formula

$$u_\delta(x) = \inf_{\gamma} \int_{-\infty}^0 e^{-\delta|t|} L(\gamma, \dot{\gamma}) dt$$

where $\gamma : (-\infty, 0] \rightarrow \mathbb{R}^d$ is absolutely continuous and $\gamma(0) = x$

- ◆ apply the Lions-Papanicolaou-Varadhan theorem (1987)

$$-\delta u_\delta \xrightarrow{C^0} \bar{H}, \quad \left\| u_\delta + \frac{\bar{H}}{\delta} \right\|_\infty \leq C, \quad \text{Lip}(u_\delta) \leq C.$$

Take a sub-sequence $\delta_i \rightarrow 0$ so that $u_\delta + \frac{\bar{H}}{\delta_i} \xrightarrow{C^0} u$

Theorem (Davini, Fathi, Iturriaga, Zavidovique, 2015)

$$u_\delta + \frac{\bar{H}}{\delta} \xrightarrow{C^0} u, \quad \text{solution of the cell equation } H(x, d_x u) = \bar{H}$$

(A true limit as $\delta \rightarrow 0$)

Discretization in time

- ◆ discretize the representation formula

$$u_\delta(x) = \inf_{\gamma} \int_{-\infty}^0 e^{-\delta|t|} L(\gamma, \dot{\gamma}) dt$$

- ◆ τ time step

$$u_{\tau,\delta}(x) = \inf_{(v_{-k})_{k \geq 0}} \sum_{k=0}^{+\infty} (1 - \tau\delta)^k \tau L(x_{-k-1}, v_{-k-1})$$

$$\exp(-\delta\tau k) \simeq (1 - \tau\delta)^k, \quad x_{-k} = x_{-k-1} + \tau v_{-k-1}$$

- ◆ Dynamical programming principle

$$u_{\tau,\delta}(x_0) = \inf_{x_{-1}} \left\{ (1 - \tau\delta) u_{\tau,\delta}(x_{-1}) + \tau L(x_{-1}, \frac{x_0 - x_{-1}}{\tau}) \right\}$$

- ◆ it is easy to see, for fixed $\delta > 0$, $u_{\tau,\delta} \rightarrow u_\delta$ as $\tau \rightarrow 0$

What happens if τ is fixed and $\delta \rightarrow 0$?

Theorem (Xifeng Su, Ph. Thieullen) $H(x, p)$ is a Tonelli Hamiltonian. Let $\tau, \delta > 0$ and $u_{\tau, \delta}$ be a solution of the discrete discounted cell equation

$$u_{\tau, \delta}(y) = \inf_{x \in \mathbb{R}^d} \left\{ (1 - \tau\delta)u_{\tau, \delta}(x) + \tau L\left(x, \frac{x - y}{\tau}\right) \right\}, \quad \forall y \in \mathbb{R}^d$$

Then as $\delta \rightarrow 0$

- ♦ uniformly in x , $\tau\delta u_{\tau, \delta}(x) \rightarrow \bar{E}_\tau$
- ♦ uniformly in x , $u_{\tau, \delta}(x) - \frac{\bar{E}_\tau}{\tau\delta} \rightarrow u_\tau(x)$, u_τ is Lipschitz
- ♦ u_τ is a particular solution of the discrete cell equation

$$u_\tau(y) + \bar{E}_\tau = \inf_{x \in \mathbb{R}^d} \left\{ u_\tau(x) + E_\tau(x, y) \right\}, \quad \forall y \in \mathbb{R}^d$$

where $E_\tau(x, y) = \tau L(x, \frac{y-x}{\tau})$ is the discrete Lagrangian

- ♦ uniformly in τ , $\|u_\tau\|_\infty \leq C$, $\text{Lip}(u_\tau) \leq C$
- ♦ for some sub-sequence $\tau \rightarrow 0$,

$$u_\tau \rightarrow u \quad \begin{array}{l} -\frac{\bar{E}_\tau}{\tau} \rightarrow \bar{H} \\ \text{(some solution of } H(x, d_x u) = \bar{H} \end{array}$$

Conclusion Discretization in time of the cell equation

\iff solving an additive eigenvalue problem of the form:

$$T_-[u] = u + \bar{E}, \quad \text{with} \quad T_-[u](y) = \inf_{x \in \mathbb{R}^d} \{u(x) + E(x, y)\}$$

($T_-[u]$ is called (backward) Lax-Oleinik operator). $E(x, y) = \tau L(x, \frac{y-x}{\tau})$

Proposition If $E(x, y)$ is C^0 , coercive $E(x, y) \rightarrow +\infty$ as $\|y - x\| \rightarrow +\infty$ and periodic $E(x + k, y + k) = E(x, y)$, then

- ♦ \bar{E} is the unique additive eigenvalue
- ♦ $\exists u$ periodic, but u may not be unique.

Remark The minimization can be taken on $[0, 1]^d$

$$\tilde{E}(x, y) = \min_{k \in \mathbb{Z}^d} E(x + k, y) \implies T_-^E[u] = T_-^{\tilde{E}}[u]$$

Program of research Identify the solution (\bar{E}, u) as the ground level of some associated dynamical system. Simplify the problem by assuming, the space is discrete, $\{x_i : i = 1 \dots N\}$ a grid of $[0, 1]^d$, and the operator is discrete

$$T_-[u](x_j) = \min_{1 \leq i \leq N} \{u(x_i) + \tilde{E}(x_i, x_j)\}$$

II. Thermodynamic formalism

Objective Give a precise meaning to the notion of configurations at equilibrium for some temperature $T > 0$.

$$\text{configuration} = (\cdots, x_{i-2}, x_{i-1} \mid x_{i_0}, x_{i_1}, \cdots), \quad x_i \in \text{grid of } [0, 1]^d$$

Notations ♦ $S = \{1, \cdots, N\}$, some finite state space

♦ $\Omega = S^{\mathbb{Z}}$, the space of all configurations,

$$\omega = (\cdots, \omega_{-1} \mid \omega_0, \omega_1, \cdots), \quad \omega_k \in S$$

♦ $\sigma : \Omega \rightarrow \Omega$, the left shift on the indices

$$\sigma(\omega) = (\cdots, \omega_{-1}, \omega_0 \mid \omega_1, \omega_2, \cdots)$$

♦ $E : \Omega \rightarrow \mathbb{R}$, some Hölder function (long range interaction)

A short range interaction: $E(\omega) = E(\omega_0, \omega_1)$

Before we had $E(x_i, x_j) = \tau L\left(x_i, \frac{x_j - x_i}{\tau}\right)$

Important notions

- stationary or invariant measure μ : a probability measure such that if

$$[i_0, \dots, i_n]_k := \{\omega \in \Omega : \omega_k = i_0, \omega_{k+1} = i_1, \dots, \omega_{k+n} = i_n\}$$

then $\mu([i_0, \dots, i_n]_k)$ is independent of k . For example

$$\text{i.i.d. measure : } \mu([i_0, \dots, i_n]) = \left(\frac{1}{N}\right)^{n+1}$$

- the entropy of an invariant measure

$$\text{Ent}_n(\mu) := \sum_{i_0 \dots i_{n-1}} -\mu([i_0, \dots, i_{n-1}]) \ln \mu([i_0, \dots, i_{n-1}])$$

$$\lim_{n \rightarrow +\infty} \frac{1}{n} \text{Ent}_n(\mu) = \text{exists} := \text{Ent}(\mu).$$

For example $\text{Ent}(\text{i.i.d.}) = \ln N$

- the free energy, let be $\beta^{-1} > 0$ called the temperature

$$\bar{E}_\beta := \inf_{\mu \text{ invariant}} \left\{ \int E d\mu - \beta^{-1} \text{Ent}(\mu) \right\}$$

Theorem (Bowen, Ruelle, ~1970) To simplify: $E(\omega) = E(\omega_0, \omega_1)$

- ◆ The infimum in the free energy

$$\bar{E}_\beta := \inf_{\mu \text{ invariant}} \left\{ \int E d\mu - \beta^{-1} \text{Ent}(\mu) \right\}$$

is reached by a unique invariant measure: called Gibbs measure at temperature β^{-1} and denoted μ_β

- ◆ There is an explicit formula

$$\mu_\beta([i_0, \dots, i_n]) = \frac{\exp[-\beta \sum_{k=0}^{n-1} E(i_k, i_{k+1})]}{Z_\beta(i_0, \dots, i_n)}$$

$$\frac{1}{Z_\beta(i_0, \dots, i_n)} = \phi_\beta^-(i_0) \exp(\beta n \bar{E}_\beta) \phi_\beta^+(i_n)$$

- ◆ ϕ_β^\pm are backward and forward eigenfunctions of the transfer operator for the largest eigenvalue $\lambda_\beta = \exp[-\beta \bar{E}_\beta]$

$$\mathcal{L}_\beta^-[\phi_\beta^-](j) := \sum_{i=1}^N \phi_\beta^-(i) \exp[-\beta E(i, j)] = \exp[-\beta \bar{E}_\beta] \phi_\beta^-(j)$$

$$\mathcal{L}_\beta^+[\phi_\beta^+](i) := \sum_{j=1}^N \exp[-\beta E(i, j)] \phi_\beta^+(j) = \exp[-\beta \bar{E}_\beta] \phi_\beta^+(i)$$

Transfer operator and Lax-Oleinik

- ◆ The backward case

$$\mathcal{L}_\beta^-[\phi](j) := \sum_{i=1}^N \phi(i) \exp[-\beta E(i, j)]$$

- ◆ The Hopf technique:

$$\phi(i) = \exp[-\beta u(i)], \quad \mathcal{L}_\beta^-[\exp(-\beta u)](j) = \exp[-\beta T_\beta^- [u](j)]$$

- ◆ let $\beta \rightarrow +\infty$ in

$$\exp[-\beta T_\beta^- [u](j)] = \sum_{i=1}^N \exp[-\beta(u(i) + E(i, j))]$$

- ◆ the zero-temperature of the transfer operator = Lax-Oleinik

$$T_\beta^- [u](j) \rightarrow T^- [u](j) := \min_{1 \leq i \leq N} [u(i) + E(i, j)]$$

Important and simple facts

◆ The eigenvalue problem for transfer operator at temperature β^{-1} and for the Lax-Oleinik operator

$$T_\beta^- [u_\beta] = u_\beta + \bar{E}_\beta \quad \text{versus} \quad T^- [u] = u + \bar{E}$$

- ◆ there exists $C > 0$, $\|u_\beta\|_\infty \leq C$, $\text{Lip}(u_\beta) \leq C$
- ◆ $\bar{E}_\beta = \inf_\mu \{ \int E d\mu - \beta^{-1} \text{Ent}(\mu) \} \rightarrow \inf_\mu \int E d\mu = \bar{E}$
- ◆ there exists $\beta_i \rightarrow +\infty$, $\mu_{\beta_i} \rightarrow \mu_{min}$ and $u_{\beta_i} \rightarrow u$

$$\mu_{min} \in \arg \min_{\mu \text{ invariant}} \int E d\mu, \quad \text{Mather}(E) = \cup_{\mu_{min}} \text{supp}(\mu_{min})$$

Conclusion In the discrete case, both in time and in space

$$H(x, d_x u) = \bar{H} \iff \begin{cases} u & \text{"="} \lim_{\beta \rightarrow +\infty} -\frac{1}{\beta} \ln \phi_\beta \\ -\bar{H} & = \lim_{\beta \rightarrow +\infty} \bar{E}_\beta = \lim_{\beta \rightarrow +\infty} -\frac{1}{\beta} \ln \lambda_\beta \end{cases}$$

- ◆ problem: the limit has to be taken along a sub-sequence

III. Selection principle

Question By freezing the system, do the Gibbs measure select a particular configuration?

Counter example (Chazottes-Hochman, 2010) There exists an Hölder energy $E : \Omega \rightarrow \mathbb{R}$, (long range), such that μ_β does not converge.

Question For which energy E do \bar{E}_β and u_β converge?

$$\begin{cases} u_\beta = -\frac{1}{\beta} \ln \phi_\beta \\ \bar{E}_\beta = -\frac{1}{\beta} \ln \lambda_\beta \end{cases} \quad \begin{cases} \mathcal{L}_\beta[\phi_\beta] = \lambda_\beta \phi_\beta \\ \mathcal{L}_\beta(i, j) = \exp[-\beta E(i, j)], \text{ (short range)} \end{cases}$$

A simpler question: for which E , do the Gibbs measure μ_β converge?

Theorem (Brémont, 2003) For short-range energy $E(i, j)$, $\mu_\beta \rightarrow \mu_\infty$ selects a particular minimizing invariant measure ($\mu_\infty \in \arg \min_\mu \int E d\mu$).

A simple example Consider $\Omega = \{1, 2\}^{\mathbb{Z}}$ and $E(i, j) = \begin{bmatrix} 0 & 100 \\ 1 & 0 \end{bmatrix}$.

The energy is short range, null at the two fixed points, and positive along any other periodic cycle

$$E(1^{\mathbb{Z}}) = 0 = E(2^{\mathbb{Z}}), \quad i_0 = i_n \Rightarrow \sum_{k=0}^{n-1} E(i_k, i_{k+1}) > 0$$

As $\mu_{\beta}([i_0, \dots, i_n]) = \exp[-\beta \sum_{k=0}^{n-1} E(i_k, i_{k+1})] / Z_{\beta}(i_0, \dots, i_n)$ the Gibbs measure chooses the configurations with the least energy

$$\mu_{\beta} \rightarrow \frac{1}{2} \delta_{1^{\mathbb{Z}}} + \frac{1}{2} \delta_{2^{\mathbb{Z}}}$$

For example $E(i, j) = \begin{bmatrix} 0.1 & 100 \\ 1 & 0 \end{bmatrix}$, $\mu_{\beta} \rightarrow \delta_{2^{\mathbb{Z}}}$

Question What can we say about the explicit convergence of μ_{β} when

$$\mathcal{L}_{\beta}(i, j) = \begin{bmatrix} 1 & e^{-111\beta} & e^{-45\beta} \\ e^{-\beta} & 1 & e^{-63\beta} \\ e^{-\beta} & e^{-\beta} & 1 \end{bmatrix}, \quad \text{or} \quad M_{\epsilon}(i, j) = \begin{bmatrix} 1 & \epsilon^a & \epsilon^b \\ \epsilon^{a'} & 1 & \epsilon^c \\ \epsilon^{b'} & \epsilon^{c'} & 1 \end{bmatrix}$$

Algorithm (E. Garibaldi, Ph. Thieullen, 2012) By induction on the dimension of the matrix $M_\epsilon(i, j)$. The framework needs to be extended

$$M_\epsilon(i, j) = A(i, j)\epsilon^{a(i, j)} + o(\epsilon^{a(i, j)})$$

$o(\epsilon^{a(i, j)}) = A_1(i, j)\epsilon^{a_1(i, j)} + A_2(i, j)\epsilon^{a_2(i, j)} + \dots =$ a Puiseux series

$$a(i, i) < a_1(i, j) < a_2(i, j) < \dots \quad (\text{but } a_k(i, j) \text{ are not rational})$$

Let $\lambda_\epsilon, \phi_\epsilon^\pm$ be the eigenvalue and eigenvectors of M_ϵ

$$\sum_i \phi_\epsilon^-(i) M_\epsilon(i, j) = \lambda_\epsilon \phi_\epsilon^-(j), \quad \sum_j M_\epsilon(i, j) \phi_\epsilon^+(j) = \lambda_\epsilon \phi_\epsilon^+(i)$$

Objective: find a Puiseux series expansion of $\lambda_\epsilon, \phi_\epsilon, \mu_\epsilon$

Observation 1 (Special case for short-range E) if $\mu_\beta \rightarrow \mu_\infty$, for some sub-sequence, then μ_∞ is a barycenter of measures supported on minimizing cycles

Observation 2 (In general) The Mather set $\cup_{\mu_{min}} \text{supp}(\mu_{min})$ may have several components; μ_∞ chooses the one with the largest topological entropy

Proposition $M_\epsilon(i, j) = A(i, j)\epsilon^{a(i, j)} + o(\epsilon^{a(i, j)}) \implies \lambda_\epsilon \sim \bar{\alpha}\epsilon^{\bar{a}}$

- Remember $\bar{E} = \inf_\mu \int E d\mu = \lim_{\beta \rightarrow +\infty} \bar{E}_\beta = \lim_{\epsilon \rightarrow 0} \frac{\ln \lambda_\epsilon}{\ln \epsilon} = \bar{a}$

Definitions

- $\bar{a} = \inf_{\text{cycles}} \frac{1}{n} \sum_{k=0}^{n-1} a(i_k, i_{k+1})$, cycle = (i_0, \dots, i_n) with $i_n = i_0$

For example

$$M_\epsilon \sim \begin{bmatrix} \epsilon^a & \epsilon^b \\ \epsilon^c & \epsilon^d \end{bmatrix}, \implies \bar{a} = \min\left(a, d, \frac{b+c}{2}\right)$$

- $G_{min} \subset S \times S$ the sub-graph of minimizing cycles.
- $A_{min} = [A(i, j)]_{(i, j) \in G_{min}}$ the restriction of A to G_{min}

For example

$$\frac{b+c}{2} < \min(a, d) \implies A_{min} = \begin{bmatrix} 0 & A(1, 2) \\ A(2, 1) & 0 \end{bmatrix}$$

- $\bar{\alpha}$ = the spectral radius of A_{min}

Reduction I Let $M_\epsilon(i, j) = A(i, j)\epsilon^{a(i, j)} + o(\epsilon^{a(i, j)})$

Then solve the discrete Lax-Oleinik equation for $a(i, j)$, that is find $u(i)$

$$\begin{cases} a(i, j) \geq u(j) - u(i) + \bar{a} & \forall i, j \\ a(i, j) = u(j) - u(i) + \bar{a} & \iff (i, j) \in G_{min} \end{cases}$$

◆ Let $\Delta_\epsilon = \text{diag}(\epsilon^{u(1)}, \dots, \epsilon^{u(N)})$, then

$$M_\epsilon^I := \Delta_\epsilon M_\epsilon \Delta_\epsilon^{-1} / \epsilon^{\bar{a}} = \begin{bmatrix} A'_{min} & 0 & 0 \\ 0 & A''_{min} & 0 \\ 0 & 0 & 0 \end{bmatrix} + N_\epsilon, \quad N_\epsilon = o(\text{Id})$$

with $\rho(A''_{min}) < \rho(A'_{min}) = \bar{a}$

$$M_\epsilon^I = \begin{bmatrix} A'_{min} & 0 \\ 0 & D \end{bmatrix} + N_\epsilon, \quad A'_{min} = \text{diag}(A^1_{min}, \dots, A^r_{min})$$

with $\rho(A^1_{min}) = \dots = \rho(A^r_{min}) = \bar{a}$

From now on M_ϵ is reduced to its normal form

Reduction II $M_\epsilon = \begin{bmatrix} A & 0 \\ 0 & D \end{bmatrix} + N_\epsilon$, $A = \text{diag}(A^1, \dots, A^r)$ with $\rho(D) < \rho(A^1) = \dots = \rho(A^r) = \bar{\alpha}$. We already know $\lambda_\epsilon \sim \bar{\alpha}$

- ◆ R_ϵ, L_ϵ the right and left eigenvectors of M_ϵ for λ_ϵ

$$\begin{cases} M_\epsilon^{AA} R_\epsilon^A + M_\epsilon^{AD} R_\epsilon^D & = \lambda_\epsilon R_\epsilon^A \\ M_\epsilon^{DA} R_\epsilon^A + M_\epsilon^{DD} R_\epsilon^D & = \lambda_\epsilon R_\epsilon^D \end{cases}$$

Extract $R_\epsilon^D = (\lambda_\epsilon - M_\epsilon^{DD})^{-1} M_\epsilon^{DA} R_\epsilon^A$ and substitute

- ◆ We are left to study M_ϵ of the form

$$M_\epsilon = \begin{bmatrix} A^1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & A^r \end{bmatrix} + N_\epsilon, \quad \rho(A^1) = \dots = \rho(A^r) = \bar{\alpha}$$

- ◆ $R_\epsilon^i, L_\epsilon^i$ the restriction of R_ϵ, L_ϵ to the indices of A^i
- ◆ R^i, L^i the right and left eigenvectors of A^i for $\bar{\alpha}$

Proposition $\frac{R_\epsilon^i(x)}{R_\epsilon^i(y)} \sim \frac{R^i(x)}{R^i(y)}, \quad \forall x \sim y$ but $\frac{R_\epsilon^i(x)}{R_\epsilon^j(y)} \rightarrow ??$ for $i \neq j$

Reduction II $M_\epsilon = \text{diag}(A^1, \dots, A^r) + N_\epsilon$, $\rho(A^1) = \dots = \rho(A^r) = \bar{\alpha}$
Write a system of r equations by blocs

$$(A^i + N_\epsilon^{ii})R_\epsilon^i + \sum_{j \neq i} N_\epsilon^{ij}R_\epsilon^j = \lambda_\epsilon R_\epsilon^i$$

Take the scalar product with the left fixed eigenvector L^i of A^i

$$\left(\frac{L^i N_\epsilon^{ii} R_\epsilon^i}{L^i R_\epsilon^i} \right) L^i R_\epsilon^i + \sum_{j \neq i} \left(\frac{L^i N_\epsilon^{ij} R_\epsilon^j}{L^i R_\epsilon^j} \right) L^i R_\epsilon^j = (\lambda_\epsilon - \bar{\alpha}) L^i R_\epsilon^i$$

$$M_\epsilon^{\text{II}}(i, j) := \left[\frac{L^i N_\epsilon^{ij} R_\epsilon^j}{L^i R_\epsilon^j} \right], \quad M_\epsilon^{\text{II}} R_\epsilon^{\text{II}} = (\lambda_\epsilon - \bar{\alpha}) R_\epsilon^{\text{II}}$$

$$M_\epsilon^{\text{II}}(i, j) = A^{\text{II}}(i, j) \epsilon^{a^{\text{II}}(i, j)} + o(\epsilon^{a^{\text{II}}(i, j)}),$$

$$a^{\text{II}}(i, j) > 0, \quad \lambda_\epsilon - \bar{\alpha} \sim \bar{\alpha}^{(2)} \epsilon^{\bar{a}^{(2)}} \quad (\text{by induction})$$

The missing case There is a problem if $r = N$

$$M_\epsilon = \bar{\alpha} \text{Id} + N_\epsilon, \quad (M_\epsilon - \bar{\alpha} \text{Id}) = N_\epsilon = B(i, j) \epsilon^{b(i, j)} + o(\epsilon^{b(i, j)})$$

$$\text{diag}(N_\epsilon) = \text{diag}(B^1) \epsilon^{b^1} + \dots + \text{diag}(B^s) \epsilon^{b^s} + N_\epsilon^{\text{II}}$$

the minimizing sub-graph of N_ϵ^{II} either contains a cycle of order at least 2, or the number of cycles of order 1 est less than the dimension of N_ϵ^{II}

An example $N = 3$, $M_\epsilon(i, j) = \exp[-\beta E(i, j)]$. Reduction I gives

- ◆ A unique dominant irreducible component: $\bar{\alpha} = 1$,

$$M_\epsilon^I = \begin{bmatrix} \epsilon^a & 1 & \epsilon^{b'} \\ \epsilon^{c'} & \epsilon^b & 1 \\ 1 & \epsilon^{a'} & \epsilon^c \end{bmatrix} \text{ or } \begin{bmatrix} \epsilon^a & 1 & \epsilon^d \\ 1 & \epsilon^b & \epsilon^e \\ \epsilon^{d'} & \epsilon^{e'} & \epsilon^c \end{bmatrix} \text{ or } \begin{bmatrix} 1 & \epsilon^a & \epsilon^c \\ \epsilon^{a'} & \epsilon^b & \epsilon^d \\ \epsilon^{c'} & \epsilon^{d'} & \epsilon^e \end{bmatrix}.$$

- ◆ Two irreducible components with equal dominant spectral radius

$$\bar{\alpha} = 1, \quad M_\epsilon^I = \begin{bmatrix} 1 & \epsilon^a & \epsilon^b \\ \epsilon^{a'} & \epsilon^c & 1 \\ \epsilon^{b'} & 1 & \epsilon^d \end{bmatrix} \text{ or } \begin{bmatrix} 1 & \epsilon^a & \epsilon^b \\ \epsilon^{a'} & 1 & \epsilon^c \\ \epsilon^{b'} & \epsilon^{c'} & \epsilon^d \end{bmatrix}.$$

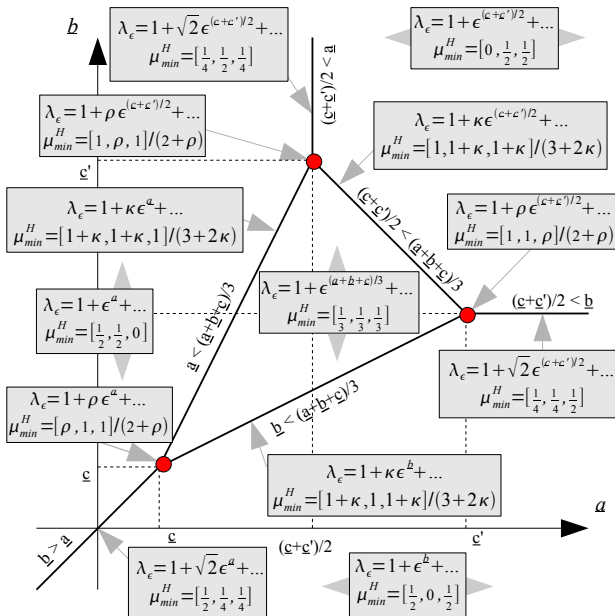
- ◆ Three irreducible components with equal dominant spectral radius:

$$\bar{\alpha} = 1, \quad M_\epsilon^I = \begin{bmatrix} 1 & \epsilon^a & \epsilon^b \\ \epsilon^{a'} & 1 & \epsilon^c \\ \epsilon^{b'} & \epsilon^{c'} & 1 \end{bmatrix}.$$






- ◆ $\mu_\beta \rightarrow$ a barycenter of the periodic minimizing cycles

$$\mu_\beta \rightarrow \mu_{min}^H := c_1 \delta_{1z} + c_2 \delta_{2z} + c_3 \delta_{3z}$$

Example A 3×3 matrix with three irreducible components



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