

Building thermodynamic formalism for hyperbolic random dynamical systems

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1. Establish **existence and uniqueness** of relative equilibrium states for random **diffeomorphism**; and
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$$\begin{aligned}\Theta : \Omega \times M &\rightarrow \Omega \times M \\ (\omega, x) &\mapsto (\sigma(\omega), T_\omega(x))\end{aligned}$$

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- $\sigma : \Omega \rightarrow \Omega$ is a homeomorphism, Ω is a compact metric space and (σ, \mathbb{P}) is ergodic.
- We also consider the projection $\pi(\omega, x) = \omega$.

A simple example

Example: Consider

$$\begin{aligned}\Theta : \Omega \times \mathbb{T}^2 &\rightarrow \Omega \times \mathbb{T}^2 \\ (\underline{\omega}, x) &\mapsto (\sigma(\underline{\omega}), Ax + \omega_0)\end{aligned}$$

- $\Omega = (\overline{B_\delta(0)})^{\mathbb{Z}}$, $\overline{B_\delta(0)}$ = closed ball of radius $\delta > 0$ centred at 0;
- $\underline{\omega} = (\omega_i)_{i \in \mathbb{Z}}$, $(\sigma \underline{\omega})_i = \omega_{i+1}$;
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- **Problem:** $h_{\text{top}}(\sigma) = \infty$ and $h_{\text{top}}(\sigma) \leq h_{\text{top}}(\Theta)$ since σ is the full shift on an uncountable alphabet.

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- **First attempt:** fix a σ -invariant \mathbb{P} and consider

$$\sup\{h_\mu(\Theta) : \pi_*\mu = \mathbb{P}\} \geq h_{\mathbb{P}}(\sigma).$$

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- **Issue:** for natural choices (e.g. $\mathbb{P} = \text{Leb}^{\otimes \mathbb{Z}}$ where Leb is the normalized Lebesgue measure on $\overline{B_\delta(0)}$), one still has $h_{\mathbb{P}}(\sigma) = \infty$.

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- **Solution:** apply a *random Brin–Katok theorem* to obtain finite local (quenched) entropy.

Relative entropy (Brin–Katok) and \mathbb{P} -relative equilibrium

Random Brin–Katok:

- $(\Theta; \sigma, \mathbb{P})$ RDS, $\Theta : \Omega \times M \rightarrow \Omega \times M$, $\Theta(\omega, x) = (\sigma\omega, T_\omega)$.
- μ is Θ -invariant, $\pi_*\mu = \mathbb{P}$ and $\mu = \mu_\omega(dx)\mathbb{P}(d\omega)$.

Recall: $\pi(\omega, x) = \omega$. Ch. 5 of B. Hasselblatt and A. Katok, eds. *Handbook of dynamical systems. Vol. 1B*. Elsevier B. V., Amsterdam, 2006, pp. xii+1222.

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Random Bowen Ball: $B_n^\omega(x, \varepsilon) := \{y : \text{dist}(T_\omega^k(y), T_\omega^k(x)) < \varepsilon, 0 \leq k < n\}$, $T_\omega^k = T_{\sigma^{k-1}\omega} \circ \cdots \circ T_\omega$.

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\mathbb{P} -relative entropy: $h_\mu(\Theta \mid \mathbb{P}) := - \iint \lim_{\varepsilon \rightarrow 0} \lim_{n \rightarrow \infty} \frac{1}{n} \log \mu_\omega(B_n^\omega(x, \varepsilon)) \, d\mu_\omega(x) \, d\mathbb{P}(\omega)$.

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Definition (\mathbb{P} -relative equilibrium state)

Let $\mathcal{M}_\mathbb{P}(\Theta) := \{\mu : \mu \text{ is } \Theta\text{-invariant and } \pi_*\mu = \mathbb{P}\}$. For measurable $\phi : \Omega \times M \rightarrow \mathbb{R} \cup \{-\infty\}$, we define the \mathbb{P} -relative topological pressure of ϕ as

$$P_{\text{top}}(\phi \mid \mathbb{P}) = \sup_{\nu \in \mathcal{M}_\mathbb{P}(\Theta)} \left\{ h_\nu(\Theta \mid \mathbb{P}) + \int \phi d\nu \right\}.$$

A maximizer $\mu \in \mathcal{M}_\mathbb{P}(\Theta)$ is a \mathbb{P} -relative equilibrium state of ϕ ; If $\phi \equiv 0$ we call μ a measure of maximal \mathbb{P} -relative entropy.

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Previous results

Random subshifts of finite type:

- Bogenschütz and Gundlach (1995)
- Mauldin and Urbański (2001)

Random countable topological Markov shifts:

Denker, Kifer, and Stadlbauer (2008)

Stadlbauer (2010,2017)

Random uniformly expanding maps:

- Kifer (1992)
- Baladi (1997)

Random non-uniformly expanding (without critical points):

- Arbieto, Matheus, and Oliveira (2004)
- Stadlbauer, Suzuki, and Varandas (2021)

Random non-uniformly hyperbolic interval maps (closed/open; discontinuities):

- Atnip, Froyland, González-Tokman, and Vaienti (2021,2023,2024,2025)

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Small perturbations of Axiom A diffeos:

- P.-D. Liu (1998)
- Chapter 5 of the book “Handbook of Dynamics” Hasselblatt and Katok (2006), written by Kifer and P.-D. Liu.

Hypothesis (H) and (H')

$$\begin{aligned}\Theta : \Omega \times M &\rightarrow \Omega \times M \\ (\omega, x) &\mapsto (\sigma\omega, T_\omega(x))\end{aligned}$$

Ω : compact metric space
 M : surface

$T_\omega : M \rightarrow M$
 $\mathcal{C}^{1+\alpha}$ diffeomorphism

(H1) (Hyperbolicity)

(H2) (Fibrewise mixing)

(H2') (Rapid fibrewise mixing)

$$(\mathbf{H}) = (\mathbf{H1}) + (\mathbf{H2}) \quad \text{and} \quad (\mathbf{H'}) = (\mathbf{H1}) + (\mathbf{H2'})$$

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(H1) (Hyperbolicity) There exist **family of deterministic cones** $\mathcal{C} = (\mathcal{C}^-(x), \mathcal{C}^+(x))_{x \in M} \subset TM$ s.t. for each $x \in M$ and \mathbb{P} -a.e. $\omega \in \Omega$,

$$DT_\omega^{-1}(x)\mathcal{C}^+(x) \subset \mathring{\mathcal{C}}^+(T_\omega^{-1}x), \quad DT_\omega(x)\mathcal{C}^-(x) \subset \mathring{\mathcal{C}}^-(T_\omega x).$$

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(H2) (Fibrewise mixing) Let $\delta > 0$, and $\{B_\delta(x_i)\}_{i=1}^k$ an open cover of M . Define

$$\mathcal{T}_{\delta, \text{mix}}(\omega) = \mathcal{T}_{\text{mix}}(\omega) := \min \{n \in \mathbb{N}; T_\omega^n(B_\delta(x_i)) \cap B_\delta(x_j) \neq \emptyset, \forall 1 \leq i, j \leq k\}.$$

For every $\delta > 0$ there exists a constant $A = A(\delta) \geq 1$, satisfying

$$\mathbb{P}[\omega \in \Omega; \mathcal{T}_{\text{mix}}(\omega) \leq A] > 0.$$

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(H2') (Rapid fibrewise mixing) For each $\delta > 0$, there exist $K = K(\delta), \kappa = \kappa(\delta) > 0$ s.t.

$$\mathbb{P}[\omega \in \Omega; \mathcal{T}_{\text{mix}}(\omega) > n] \leq Ke^{-\kappa n} \text{ for every } n \in \mathbb{N}.$$

$$(\mathbf{H}) = (\mathbf{H1}) + (\mathbf{H2}) \quad \text{and} \quad (\mathbf{H'}) = (\mathbf{H1}) + (\mathbf{H2'})$$

Examples

The two examples below satisfy Hypothesis **(H')**

(1) Let $A_1, \dots, A_k \in \text{SL}(2, \mathbb{N})$ and $\text{Tr}(A_i) > 2$ for $i \in \{1, \dots, k\}$.

$$\begin{aligned}\Theta : \{1, \dots, k\}^{\mathbb{Z}} \times \mathbb{T}^2 &\rightarrow \{1, \dots, k\}^{\mathbb{Z}} \times \mathbb{T}^2 \\ (\underline{\omega}, x) &\mapsto (\sigma \underline{\omega}, A_{\omega_0}(x))\end{aligned}$$

where σ is the left shift and \mathbb{P} is a Gibbs state of σ .

(2) Let T_1, \dots, T_k be Anosov diffeomorphisms in \mathbb{T}^2 preserving the same deterministic family of cones in $T\mathbb{T}^2$

$$\begin{aligned}\Theta : \{1, \dots, k\}^{\mathbb{Z}} \times \mathbb{T}^2 &\rightarrow \{1, \dots, k\}^{\mathbb{Z}} \times \mathbb{T}^2 \\ (\underline{\omega}, x) &\mapsto (\sigma \underline{\omega}, T_{\omega_0}(x))\end{aligned}$$

where σ is the left shift and \mathbb{P} is Bernoulli.

Theorem 1 (Amorim-C.-Sassoul-Vaienti, 2025+)

Let $(\Theta; \sigma, \mathbb{P})$ be a random dynamical system satisfying **(H)** and $\phi : \Omega \times M \rightarrow \mathbb{R}$ be a random potential such that

$$\operatorname{ess\,sup}_{\omega \in \Omega} \|\phi(\omega, \cdot)\|_{C^\alpha} < \infty$$

for some $\alpha > 0$. Then, *there exists a unique \mathbb{P} -relative equilibrium state* $\nu \in \mathcal{M}_{\mathbb{P}}(\Theta)$ of ϕ .

Theorem 2 (Amorim-C.-Saussol-Vaienti, 2025+)

Let $(\Theta; \sigma, \mathbb{P})$ be a random dynamical system satisfying **(H')**, and $\phi : \Omega \times M \rightarrow \mathbb{R}$ as in Theorem 1. Then $\nu(d\omega, dx) = \nu_\omega(dx)\mathbb{P}(d\omega)$ exhibits *quenched exponential decay of correlations* for Hölder observables: For every $p \in [1, \infty)$, $\exists C_p \in L^p(\Omega, \mathbb{P})$, $\alpha > 0$ such that for any $f, g \in C^\alpha$, $\forall n \in \mathbb{N}$

$$\left| \int_M f \circ T_\omega^n \cdot g \, d\nu_\omega - \int_M f \, d\nu_{\sigma^n \omega} \int_M g \, d\nu_\omega \right| \leq C_p(\omega) e^{-\alpha_0 n} \|f\|_{C^\alpha} \|g\|_{C^\alpha},$$

for \mathbb{P} -a.e. $\omega \in \Omega$

Results

Theorem 3 (Quenched CLT)

Let $(\Theta; \sigma, \mathbb{P})$ satisfy **(H')**, let $\phi : \Omega \times M \rightarrow \mathbb{R}$ be as in Theorem 1, and let $\nu = \nu_\omega \mathbb{P}(d\omega)$ be the \mathbb{P} -relative equilibrium state. Then ν satisfies the **quenched CLT** for any observable $f : \Omega \times M \rightarrow \mathbb{R}$ such that

$$\operatorname{ess\,sup}_{\omega \in \Omega} \|f(\omega, \cdot)\|_{C^\alpha} < \infty.$$

Set

$$\tilde{f}(\omega, x) := f(\omega, x) - \int_M f(\omega, \cdot) d\nu_\omega, \quad S_n^\omega \tilde{f}(x) := \sum_{k=0}^{n-1} \tilde{f}(\sigma^k \omega, T_\omega^k(x)).$$

There exists $\Sigma^2 \geq 0$, independent of ω , such that for \mathbb{P} -a.e. ω ,

$$\frac{1}{\sqrt{n}} S_n^\omega \tilde{f} \xrightarrow{n \rightarrow \infty} \mathcal{N}(0, \Sigma^2),$$

Moreover, \tilde{f} is not a **quenched coboundary**, i.e.,

it **does not** exist $\psi \in L^2(\mathbb{P} \otimes \nu)$ such that $\tilde{f} = g \circ \Theta - g$ for \mathbb{P} -a.e. ω , ν_ω -a.e. x .

then $\Sigma > 0$.

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- (i) **Random Banach spaces:** One finds an appropriate family of Banach spaces $(\mathcal{B}_\omega)_{\omega \in \Omega}$ such that \mathcal{C}^α is dense in \mathcal{B}_ω and

$$\mathcal{L}_\omega : \mathcal{B}_\omega \rightarrow \mathcal{B}_{\sigma\omega}, \quad \mathcal{L}_\omega(f) := (e^{\phi_\omega} f) \circ T_\omega^{-1} \quad (\forall f \in \mathcal{C}^\alpha),$$

is bounded, where $\phi_\omega(\cdot) := \phi(\omega, \cdot)$. We also ask $\|\cdot\|_{\mathcal{B}_\omega} \leq \|\cdot\|_{\mathcal{C}^\alpha}$

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- (ii) **Random cones:** For \mathbb{P} -a.e. $\omega \in \Omega$, one finds an appropriate family of cones $\mathcal{C}_\omega \subset \mathcal{B}_\omega$, i.e.

$$\mathcal{C}_\omega \cap (-\mathcal{C}_\omega) = \{0\}; \quad f \in \mathcal{C}_\omega, \quad \alpha > 0 \Rightarrow \alpha f \in \mathcal{C}_\omega; \quad f, g \in \mathcal{C}_\omega \Rightarrow f + g \in \mathcal{C}_\omega,$$

such that $\mathcal{L}_\omega(\mathcal{C}_\omega) \subset \mathcal{C}_{\sigma\omega}$ and $\text{span}\{\mathcal{C}_\omega\} = \mathcal{B}_\omega$.

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$$\mathcal{C}_\omega \cap (-\mathcal{C}_\omega) = \{0\}; \quad f \in \mathcal{C}_\omega, \quad \alpha > 0 \Rightarrow \alpha f \in \mathcal{C}_\omega; \quad f, g \in \mathcal{C}_\omega \Rightarrow f + g \in \mathcal{C}_\omega,$$

such that $\mathcal{L}_\omega(\mathcal{C}_\omega) \subset \mathcal{C}_{\sigma\omega}$ and $\text{span}\{\mathcal{C}_\omega\} = \mathcal{B}_\omega$.

- (iii) **Cone contraction**: For \mathbb{P} -a.e. $\omega \in \Omega$ there exist random variables

$$\{\tau_i : \Omega \rightarrow \mathbb{N}\}_{i \in \mathbb{N}} \text{ with } \tau_1 \leq \tau_2 \leq \tau_3 \leq \dots$$

and constants $\Delta > 0$ and $\chi \in (0, 1)$ (independent of ω) such that, for all $k \geq 1$,

$$\sup_{f, g \in \mathcal{C}_{\sigma^{-\tau_k}\omega}} \text{Hil}_{\mathcal{C}_\omega} \left(\mathcal{L}_{\sigma^{-\tau_k}\omega}^{(\tau_k)} f, \mathcal{L}_{\sigma^{-\tau_k}\omega}^{(\tau_k)} g \right) \leq \Delta \chi^{k-1}; \quad \left(\text{recall that } \mathcal{L}_{\sigma^{-\tau_k}\omega}^{(\tau_k)} \mathcal{C}_{\sigma^{-\tau_k}\omega} \subset \mathcal{C}_\omega \right),$$

where $\text{Hil}_{\mathcal{C}_\xi}$ is the **projective Hilbert metric** on \mathcal{C}_ξ and $\mathcal{L}_\xi^{(n)} := \mathcal{L}_{\sigma^{n-1}\xi} \circ \dots \circ \mathcal{L}_\xi$.

For a cone \mathcal{C} with order $\leq_{\mathcal{C}}$, set $\alpha(f, g) := \sup\{\lambda > 0 : \lambda g \leq_{\mathcal{C}} f\}$, $\beta(f, g) := \inf\{\mu > 0 : f \leq_{\mathcal{C}} \mu g\}$, and $\text{Hil}_{\mathcal{C}}(f, g) := \log \frac{\beta(f, g)}{\alpha(f, g)}$.

Right and left eigenvectors, and spectral gap

Assume cones $\mathcal{C}_\omega \subset \mathcal{B}_\omega$ and dual cones $\mathcal{C}_\omega^* \subset \mathcal{B}_\omega^*$ and that (i)–(iii) from the previous slide hold. Then via standard arguments, for \mathbb{P} -a.e. ω we have that:

Adaptation of the Birkhoff cones arguments to the Liverani, Saussol, and Vaienti (1998) to the random context

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Right random eigenvector: recall that $\mathcal{L}_\omega : \mathcal{C}_\omega \rightarrow \mathcal{C}_{\sigma\omega}$, there exists $\mu_\omega \in \mathcal{C}_\omega$ and $\lambda_\omega > 0$ satisfying

$$\mu_\omega := \lim_{n \rightarrow \infty} \frac{\mathcal{L}_{\sigma^{-n}\omega}^n \mathbb{1}}{\|\mathcal{L}_{\sigma^{-n}\omega}^{(n)} \mathbb{1}\|_{\mathcal{B}_\omega}} \in \mathcal{C}_\omega, \quad \mathcal{L}_\omega \mu_\omega = \lambda_\omega \mu_{\sigma\omega}.$$

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Spectral gap:

$$\mathcal{L}_\omega^{(n)} g = \lambda_\omega^{(n)} \ell_\omega(g) \mu_{\sigma^n \omega} + \lambda_\omega^{(n)} Q_\omega^{(n)}(g), \quad \|Q_\omega^{(n)}\|_{\mathcal{B}_\omega \rightarrow \mathcal{B}_{\sigma^n \omega}} \leq C \chi^{N_\tau(\omega, n)},$$

where

$$\lambda_\omega^{(n)} := \prod_{j=0}^{n-1} \lambda_{\sigma^j \omega} \quad \text{and} \quad N_\tau(\omega, n) = \#(\{\tau_i(\sigma^n \omega)\}_{i \in \mathbb{N}} \cap [1, n])$$

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Natural candidates

For each $f \in \mathcal{C}^\alpha(M)$, define $\nu_\omega(f) = \ell_\omega(f \cdot \mu_\omega)$. Observe that ν_ω is **equivariant**:

$$\nu_\omega(f \circ T_\omega) = \ell_\omega(f \circ T_\omega \cdot \mu_\omega) = \frac{1}{\lambda_\omega}(\mathcal{L}_\omega^*) \ell_{\sigma\omega}(f \circ T_\omega \cdot \mu_\omega) = \frac{1}{\lambda_\omega} \ell_\omega(f \cdot \mathcal{L}_\omega \mu_\omega) = \nu_{\sigma\omega}(f).$$

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Also, choosing $f, g \in \mathcal{C}^\alpha(M)$ we have

$$\begin{aligned} \nu_\omega(f \circ T_\omega^n \cdot g) &= \ell_\omega(f \circ T_\omega^n \cdot g \cdot \mu_\omega) = \frac{1}{\lambda_\omega^{(n)}} (\mathcal{L}_\omega^*)^{(n)} \ell_{\sigma^n\omega}(f \circ T_\omega^n \cdot g \cdot \mu_\omega) = \frac{1}{\lambda_\omega^{(n)}} \ell_{\sigma^n\omega}\left(f \cdot \mathcal{L}_\omega^{(n)}(g \cdot \mu_{\sigma^n\omega})\right) \\ &= \ell_{\sigma^n\omega}\left(f \cdot \ell_\omega(g \cdot \mu_\omega) \mu_\omega + f \cdot Q_\omega^{(n)}(g \mu_\omega)\right) = \nu_{\sigma^n\omega}(f) \nu_\omega(g) + \mathcal{O}\left(\|Q_\omega^{(n)}\|_{B_\omega \rightarrow B_{\sigma^n\omega}}\right). \end{aligned}$$

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If one manages to do such a construction:

- $\nu = \nu_\omega(dx) \mathbb{P}(d\omega)$ is a **candidate for \mathbb{P} -relative equilibrium state** for the potential

$$\phi + \log \left| \det \left(D T_\omega|_{E^s(\omega, x)} \right) \right|.$$

- $\int \log \lambda_\omega \mathbb{P}(d\omega)$ is a **candidate for its \mathbb{P} -relative topological pressure**.

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Main challenge: There is *no thermodynamic formalism for diffeomorphisms based on Birkhoff cones*. Existing results (Baladi–Tsujii, 2007; Gouëzel–Liverani, 2008) rely on spectral methods (anisotropic Banach spaces) instead.

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(C3) Hölder across leaves (holonomy): for holonomy-related $(\tilde{\gamma}_\omega, \gamma_\omega)$, if $\tilde{\rho}_\omega$ is the transport/normalisation of ρ_ω ,

$$\frac{\int_{\tilde{\gamma}_\omega} \varphi \tilde{\rho}_\omega \, d\mathbf{m}_{\tilde{\gamma}_\omega}}{\int_{\gamma_\omega} \varphi \rho_\omega \, d\mathbf{m}_{\gamma_\omega}} \leq e^{c(d_u(\tilde{\gamma}_\omega, \gamma_\omega))^\eta}$$

Cone definition

We define \mathcal{C}_ω as the closure of $C_\omega \subset \{\varphi : M \rightarrow \mathbb{R}; \varphi \text{ bounded and measurable}\}$ with respect to the norm

$$\|f\|_\omega := \|f\|_{\omega, a, \kappa}^{\sup_s} + \frac{1}{b} \|f\|_{\omega, a, \kappa}^{\Theta_s} + \frac{1}{c} \|f\|_{\omega, \eta}^{d_u},$$

where

$$\|f\|_{\omega, a, \kappa}^{\sup_s} := \sup_{\gamma_\omega \in \Gamma_\delta^s(\omega)} \sup_{\substack{\rho_\omega \in \mathcal{D}_\omega(\gamma_\omega) \\ \int_{\gamma_\omega} \rho_\omega \, dm_{\gamma_\omega} = 1}} \left| \int_{\gamma_\omega} f \rho \, dm_{\gamma_\omega} \right|,$$

$$\|f\|_{\omega, a, \kappa}^{\Theta_s} := \sup_{\gamma_\omega \in \Gamma_\delta^s(\omega)} \sup_{\substack{\rho_\omega^1, \rho_\omega^2 \in \mathcal{D}_\omega(\gamma_\omega) \\ \int_{\gamma_\omega} \rho_\omega^1 \, dm_{\gamma_\omega} = 1 \\ \int_{\gamma_\omega} \rho_\omega^2 \, dm_{\gamma_\omega} = 1}} \frac{\left| \int_{\gamma_\omega} f \rho_\omega^1 \, dm_{\gamma_\omega} - \int_{\gamma_\omega} f \rho_\omega^2 \, dm_{\gamma_\omega} \right|}{\text{Hil}_{\mathcal{D}_\omega(\gamma_\omega)}(\rho_\omega^1, \rho_\omega^2)},$$

$$\|f\|_{\omega, \nu, a, \kappa_1}^{d_u} := \sup_{\substack{(\gamma_\omega, \tilde{\gamma}_\omega) \in \Gamma_\delta^s(\omega) \times \Gamma_\delta^s(\omega) \\ \text{nearby pair}}} \sup_{\rho \in \mathcal{D}_\omega(\gamma_\omega)} \frac{\left| \int_{\gamma_\omega} f \rho_\omega \, dm_{\gamma_\omega} - \int_{\tilde{\gamma}_\omega} f \tilde{\rho}_\omega \, dm_{\tilde{\gamma}_\omega} \right|}{d_u(\gamma_\omega, \tilde{\gamma}_\omega)^\eta}.$$

And consider $\mathcal{B}_\omega := \text{span}(\mathcal{C}_\omega)$.

The general strategy required us to

(i) Construct random Banach spaces \mathcal{B}_ω , (ii) Construct random cones \mathcal{C}_ω

(iii) Construct stopping times that generate cone contractions, i.e. existence of random variables

$$\tau_1 \leq \tau_2 \leq \dots : \Omega \rightarrow \mathbb{N}, \text{ satisfying } \text{diam}_{\text{Hil}_{\mathcal{C}_\omega}} \left(\mathcal{L}_{\theta^{-\tau_m}\omega}^{(\tau_m)} \mathcal{C}_{\theta^{-\tau_m}\omega} \right) \leq \Delta \chi^{m-1}.$$

Using **(H1)** (hyperbolicity), we have established (i)+(ii). We now use **(H2)** (fibrewise mixing) to obtain (iii).

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- (i) Construct random Banach spaces \mathcal{B}_ω ,
- (ii) Construct random cones \mathcal{C}_ω
- (iii) Construct stopping times that generate cone contractions, i.e. existence of random variables $\tau_1 \leq \tau_2 \leq \dots : \Omega \rightarrow \mathbb{N}$, satisfying $\text{diam}_{\text{Hil}_{\mathcal{C}_\omega}} \left(\mathcal{L}_{\theta^{-\tau_m}\omega}^{(\tau_m)} \mathcal{C}_{\theta^{-\tau_m}\omega} \right) \leq \Delta \chi^{m-1}$.

Using **(H1)** (hyperbolicity), we have established (i)+(ii). We now use **(H2)** (fibrewise mixing) to obtain (iii).

Lemma 4

Recall (from the mixing condition H2) that

$$\mathcal{T}_{\text{mix}}(\omega) := \min \left\{ n \in \mathbb{N}; T_\omega^n(B_{\delta/2}(x_i)) \cap B_{\delta/2}(x_j) \neq \emptyset, \forall 1 \leq i, j \leq k \right\}.$$

Then there exists $D_1, D_2 > 1$ independent of ω such that for any $n \geq \mathcal{T}_{\text{mix}}(\omega)$

$$\sup_{f, g \in \mathcal{C}} \text{Hil}_{\mathcal{C}_\omega} \left(\mathcal{L}_\omega^{(n)} f, \mathcal{L}_\omega^{(n)} g \right) \leq D_1 + D_2^{\mathcal{T}_{\text{mix}}(\omega)}.$$

In particular if $\mathcal{T}_{\text{mix}}(\omega) \leq A$, then $\sup_{f, g \in \mathcal{C}_\omega} \text{Hil}_{\mathcal{C}_\omega} \left(\mathcal{L}_\omega^{(n)} f, \mathcal{L}_\omega^{(n)} g \right) \leq D_1 + D_2^A := \Delta$.

See Liverani (1995) or X. Liu (2024) (see also Atnip–Froyland–Gonzalez-Tokman–Vaienti (2021) and Buzzi (1999)).

(iii) Construction of stopping times that generate cone contractions

By **(H2)** there exists A such that $\mathbb{P}[\mathcal{T}_{\text{mix}}(\omega) \leq A] > 0$.

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$$\tau_0 = 0 < \tau_1 < \tau_2 < \cdots \text{ with } N(\sigma^{-\tau_m}\omega) \leq A \text{ and } \tau_m - \tau_{m-1} \geq A \text{ for every } m \in \mathbb{N}.$$

Let Δ as before and $\chi := \tanh(\Delta/4) \in (0, 1)$. Then for any $m \in \mathbb{N}$,

$$\sup_{f, g \in \mathcal{C}_{\sigma^{-\tau_m}\omega}} \text{Hil}_{\mathcal{C}_{\sigma^{-\tau_m}\omega}} \left(\mathcal{L}_{\sigma^{-\tau_m}\omega}^{(\tau_m)} f, \mathcal{L}_{\sigma^{-\tau_m}\omega}^{(\tau_m)} g \right) \leq \chi^{m-1} \Delta.$$

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Consequence of the previous lemma and

$$\mathcal{L}_{\sigma^{-\tau_m}\omega}^{(\tau_m)} \mathcal{C}_{\sigma^{-\tau_m}\omega} = \mathcal{L}_{\sigma^{-\tau_1}\omega}^{(\tau_1)} \circ \mathcal{L}_{\sigma^{-\tau_2}\omega}^{(\tau_2 - \tau_1)} \circ \cdots \circ \mathcal{L}_{\sigma^{-\tau_m}\omega}^{(\tau_m - \tau_{m-1})} \mathcal{C}_{\sigma^{-\tau_m}\omega} \subset \mathcal{C}_\omega,$$

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From the discussion in the previous slides, we can construct

- $\nu_\omega(f) := \ell_\omega(f \mu_\omega)$, with $\mathcal{L}_\omega \mu_\omega = \lambda_\omega \mu_{\sigma\omega}$ and $(\mathcal{L}_\omega)^* \ell_{\sigma\omega} = \lambda_\omega \ell_\omega$;
- $\left| \int_M f \circ T_\omega^n g \, d\nu_\omega - \int f \, d\nu_{\sigma^n\omega} \int_M g \, d\nu_\omega \right| \leq \chi^{N_\tau(\omega, n)} \|f\|_{C^\alpha} \|g\|_{C^\alpha}$, with

$$N_\tau(\omega, n) := \#(\{\tau_i(\sigma^n\omega)\}_{i \in \mathbb{N}} \cap [1, n]).$$

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In particular, if **(H2')** is satisfied ($\mathbb{P}[\mathcal{T}_{\text{mix}}(\omega) > n] \leq C_0 e^{-\kappa n}$), then the decay is exponentially fast (depending on ω).

Showing that ν is the unique \mathbb{P} -relative equilibrium state

Assume that the RDS $(\Theta; \sigma, \mathbb{P})$ satisfies Hypothesis **(H)**. Let $\nu = \nu_\omega(\cdot) \mathbb{P}(d\omega)$, with $\nu_\omega(f) = \ell_\omega(f \cdot \mu_\omega)$.

- We show that

$$\int \log \lambda_\omega \, d\mathbb{P}(\omega) = P_{\text{top}}(\phi + \log |\det(DT_\omega|_{E^s(\omega, x)})| \mid \mathbb{P}),$$

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- We show uniqueness of equilibrium states by noting that ℓ_ω and μ_ω are Margulis measures and following the method described in Section 4 of Carrasco and Rodríguez-Hertz (2023).
- We prove quenched CLT by the standard Nagaev-Guivarc'h perturbation method for complex cones, Chapter 7 of Hafouta, and Kifer (2018).

Results

- We have established **existence and uniqueness of relative equilibrium states** for hyperbolic random dynamical systems on surfaces, under the uniform Hölder bound

$$\operatorname{ess\,sup}_{\omega} \|\phi(\omega, \cdot)\|_{\mathcal{C}^{\alpha}} < \infty.$$

- We have obtained **quenched exponential decay of correlations** and a **quenched CLT** via the basis large deviation estimates, i.e.

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Thank you for your attention!

Topological pressure via the random eigenvalues

Let $\nu = \nu_\omega(dx)\mathbb{P}(d\omega)$ and assume $(\mathbf{H}) = (\mathbf{H}_1) + (\mathbf{H}_2)$ (Hyperbolicity + Fibrewise mixing).

- We show that

$$\int \log \lambda_\omega \, d\mathbb{P}(\omega) = P_{\text{top}}(\Theta, \phi(\omega, x) + \log |\det(DT_\omega|_{E^s(\omega, x)})| \mid \mathbb{P}),$$

by adapting an argument of Parmenter–Pollicott (2021) which reads

Theorem 5 (Parmenter–Pollicott (2021))—

$T : M \rightarrow M$ is a mixing Anosov map (or has an Axiom A attractor) and $\phi : M \rightarrow \mathbb{R}$ is continuous, then for any piece of stable manifold $\gamma \subset M$,

$$P(\phi + \log |\det(DT|_{E^s})|, T) = \lim_{n \rightarrow \infty} \frac{1}{n} \log \int_\gamma \exp\left(\sum_{i=0}^{n-1} \phi \circ T^{-i}(y)\right) dm_\gamma(y),$$

where m_γ is the induced Riemannian measure on γ .

The measure ν is a \mathbb{P} -relative equilibrium state

Let $\nu = \nu_\omega(dx)\mathbb{P}(d\omega)$ and assume $(\mathbf{H}) = (\mathbf{H}_1) + (\mathbf{H}_2)$ (Hyperbolicity + Fibrewise mixing).

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$$\int \log \lambda_\omega d\mathbb{P}(\omega) = P_\nu(\Theta, \phi(\omega, x) + \log |\det(DT_\omega|_{E^s(\omega, x)})| \mid \mathbb{P}),$$

By establishing a weak Gibbs property:

Proposition 1 (Weak Gibbs property)

Let $\varepsilon > 0$ be small enough. Then, there exist functions $K_\varepsilon \in L^1(\Omega, \mathbb{P})$ and measurable functions $c_\varepsilon, C_\varepsilon : \Omega \rightarrow (0, \infty)$, such that for \mathbb{P} -a.e. $\omega \in \Omega$ there exists a strictly increasing sequence $\{n_k(\omega)\}_{k \in \mathbb{N}}$ such that for every $k \in \mathbb{N}$ and $x \in M$:

$$c_\varepsilon(\omega) e^{-K_\varepsilon(\sigma^{-n_k(\omega)}\omega)} \leq \frac{\nu_\omega(B_\omega^{n_k(\omega)}(x, \varepsilon))}{[\lambda_{\sigma^{-n_k(\omega)}\omega}^{(n_k(\omega))}]^{-1} \exp\left(S_n \bar{\phi}_{\sigma^{-n_k(\omega)}\omega}\left((T_{\sigma^{-n_k(\omega)}\omega}^n)^{-1}x\right)\right)} \leq C_\varepsilon(\omega) e^{K_\varepsilon(\sigma^{-n_k(\omega)}\omega)},$$

where $B_\omega^n(x, \varepsilon) = \{y \in M : \forall 0 \leq i \leq n-1, d(T_\omega^{-i}x, T_\omega^{-i}y) \leq \varepsilon\}$ is the backward dynamical ball.

The proof of the above proposition borrows some ideas from Stadlbauer–Suzuki–Varandas (2021).

Uniqueness of \mathbb{P} -relative equilibrium states

Let $\nu = \nu_\omega(dx) \mathbb{P}(d\omega)$ and assume $(\mathbf{H}) = (\mathbf{H}_1) + (\mathbf{H}_2)$ (hyperbolicity + fibrewise mixing).

- **Deterministic result (Carrasco–Rodríguez–Hertz, 2023).** If $T : M \rightarrow M$ is a centre isometry (i.e. $TM = E^s \oplus E^c \oplus E^u$ with E^c having special properties) and $\bar{\phi}$ is a suitable potential (in our context $E^c = \emptyset$ and $\bar{\phi}$ is Hölder), and if there exist measures μ^u, μ^s such that

- (i) the (Rokhlin) disintegration of μ^u along unstable leaves $\{\mu_x^u\}$ is absolutely continuous with respect to the leafwise Lebesgue measure;
- (ii) the disintegration of μ^s along stable leaves $\{\mu_x^s\}$ is absolutely continuous with respect to the leafwise Lebesgue measure;
- (iii) *Margulis property*: for every $x \in M$,

$$T_*\mu_x^u = e^{-P_{\text{top}}(\bar{\phi}) + \bar{\phi}} \mu_{T(x)}^u, \quad T_*\mu_x^s = e^{P_{\text{top}}(\bar{\phi}) - \bar{\phi}} \mu_{T(x)}^s;$$

- (iv) T is topologically mixing,
then T admits a unique equilibrium state for $\bar{\phi}$.

- **Random adaptation.** We adapt the above result by observing that ℓ_ω and μ_ω are Margulis measures for the potential

$$\bar{\phi}(\omega, x) := \phi(\omega, x) + \log |\det(DT_\omega|_{E^s(\omega, x)})|$$

and that both μ_ω and ℓ_ω have full support.

Quenched CLT

Let $\nu = \nu_\omega(dx) \mathbb{P}(d\omega)$ and assume $(\mathbf{H}') = (\mathbf{H}_1) + (\mathbf{H}_2')$ (hyperbolicity + fibrewise rapid mixing). Take

$$f_\omega \in L^\infty(\Omega, C^\alpha(M)) \text{ with } \nu_\omega(f_\omega) = 0 \text{ and } S_n f(\omega, x) := \sum_{k=0}^{n-1} f_{\sigma^k \omega}(T_\omega^k x).$$

Normalise, without loss of generality, to $\lambda_\omega \equiv 1$ by $\bar{\phi}_\omega := \phi_\omega - \log \lambda_\omega$.

Nagaev–Guivarc’h on complex cones. Consider the twist operator

$$\mathcal{L}_{\omega,t} g := \mathcal{L}_\omega(e^{itf_\omega} g), \quad |t| \leq t_0.$$

Cone contraction (complex Hilbert metric) gives a uniform spectral gap for small $|t|$:

$$\mathcal{L}_{\omega,t}^{(n)} = \lambda_\omega(t)^{(n)} \Pi_{\sigma^n \omega}(t) + Q_{\omega,t}^{(n)}, \quad \|Q_{\omega,t}^{(n)}\| \leq C e^{-\chi n},$$

with $t \mapsto \lambda_\omega(t)$ analytic at 0, $\lambda'_\omega(0) = 0$, $\lambda''_\omega(0) = \Sigma^2$. Hence $\log \nu_\omega(e^{itS_n f}) = \frac{1}{2} \sigma^2 t^2 n + o(nt^2)$, and Lévy’s theorem completes the proof.

- Chapter 7 of [Y. Hafouta and Y. Kifer](#). *Nonconventional limit theorems and random dynamics*. World Scientific Publishing Co. Pte. Ltd., Hackensack, NJ, 2018, pp. xiii+284
- [J. Atnip, G. Froyland, C. González-Tokman, and S. Vaienti](#). “Thermodynamic Formalism and Perturbation Formulae for Quenched Random Open Dynamical Systems”. In: *Dissertationes Mathematicae* (2024). to appear; see [arXiv:2307.00774](#). [arXiv: 2307.00774 \[math.DS\]](#)
- Sequential Billiards: [M. F. Demers and C. Liverani](#). *Central Limit Theorem for Sequential Dynamical Systems*. 2025. [arXiv: 2502.07765 \[math.DS\]](#)

The log $|\det dT|_{E^s}|$ correction

We focus on the deterministic case. Take $\phi = -\log |\det dT|$, we know for this choice of potential the eigenvectors of $\mathcal{L}_\phi f = (e^\phi f) \circ T^{-1}$ should give rise to the SRB measure.

From Pesin's formula

$$P_{\text{top}}(-\log |\det dT|_{E^u}|) = h_{\mu_{\text{SRB}}} - \int \log |\det dT|_{E^u}| d\mu_{\text{SRB}} = 0.$$

we observe that a correction is needed.

Angle identity:

$$|\det dT_x| = \frac{\alpha(Tx)}{\alpha(x)} |\det dT_x|_{E^s}| |\det dT_x|_{E^u}|, \quad \alpha(x) = \sin \angle(E_x^s, E_x^u).$$

Coboundary:

$$-\log |\det dT_x| + \log |\det dT_x|_{E^s}| = -\log |\det dT_x|_{E^u}| + \log \alpha(Tx) - \log \alpha(x).$$

Leafwise change of variables (stable leaf γ):

$$\int_\gamma \mathcal{L}_\phi f \, dm_\gamma = \int_\gamma (e^\phi f) \circ T^{-1} \, dm_\gamma = \int_{T^{-1}\gamma} e^\phi f |\det dT|_{E^s}| \, dm_{T^{-1}\gamma} = \int_{T^{-1}\gamma} e^{\phi + \log |\det dT|_{E^s}|} f \, dm_{T^{-1}\gamma}.$$